

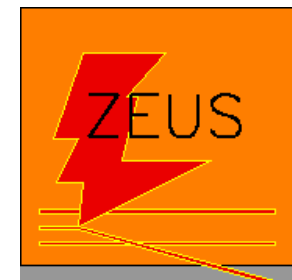
# QCD at HERA

Nick Brook



University of Bristol

- Measurement of  $F_2$
- BFKL dynamics
- Event Shapes



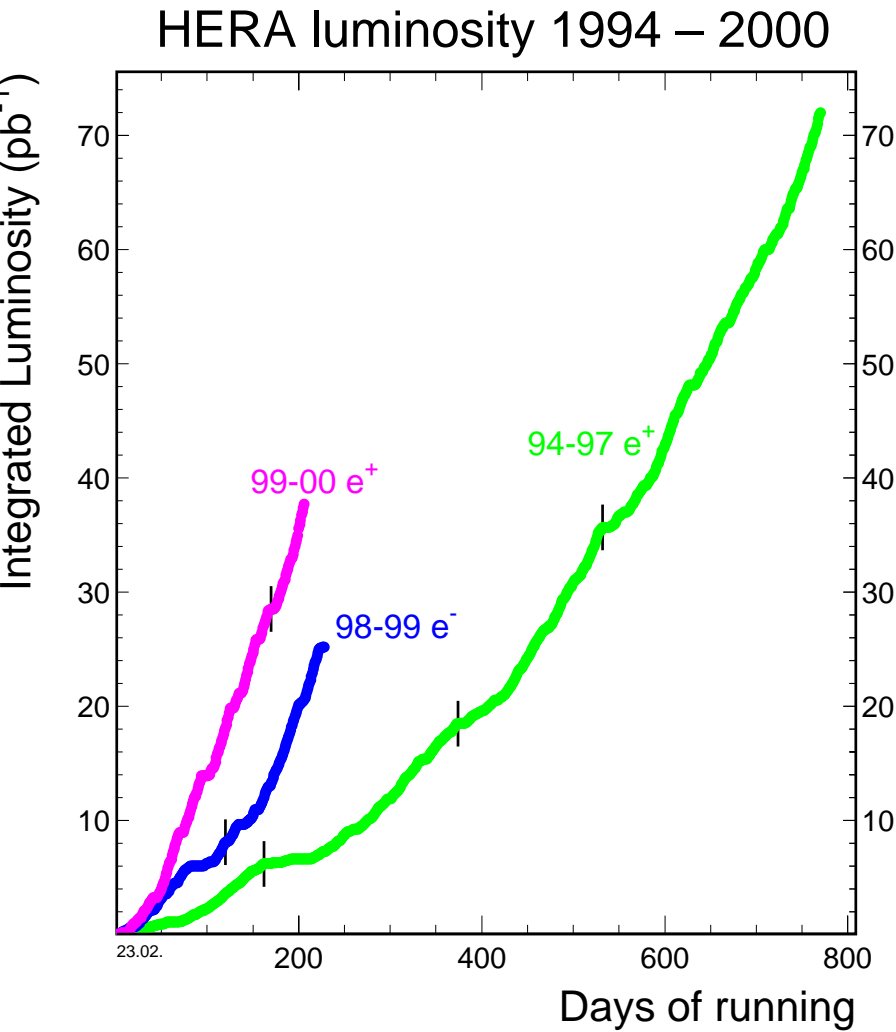
# HERA Accelerator



$e^+/e^-$  beam - 27.5 GeV

Proton beam - 820/920 GeV

# Luminosity available for physics (ZEUS):



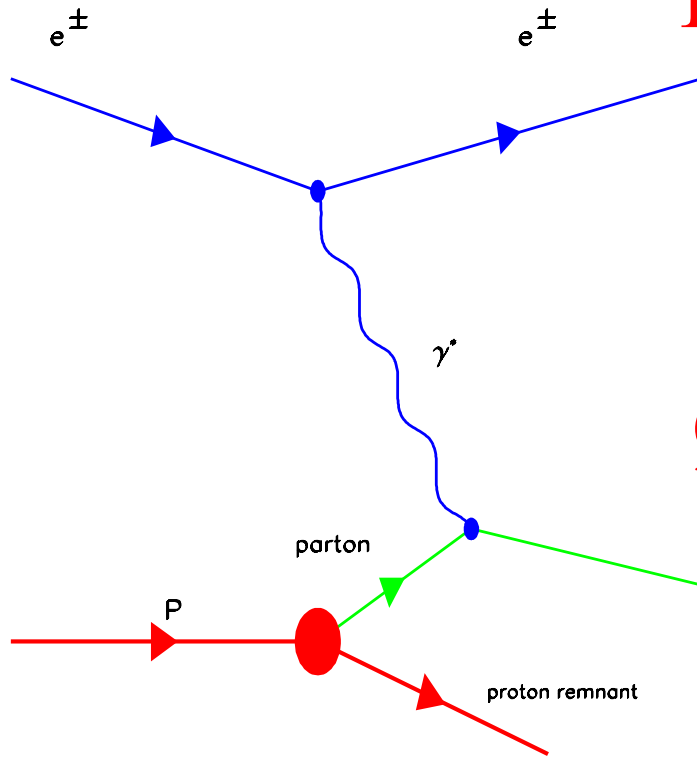
1994-97  $\sim 48 \text{ pb}^{-1}$   $e^+p$   
 $\sqrt{s} = 300 \text{ GeV}$

1998-99  $\sim 17 \text{ pb}^{-1}$   $e^-p$   
 $\sqrt{s} = 320 \text{ GeV}$

1999-2000  $\sim 25 \text{ pb}^{-1}$   $e^+p$   
 $\sqrt{s} = 320 \text{ GeV}$

**Total**  $\sim 90 \text{ pb}^{-1}$   
+ running until  
September 2000

# Naïve Quark Parton Model (QPM)



Lorentz scalars:  $y, Q^2, x$

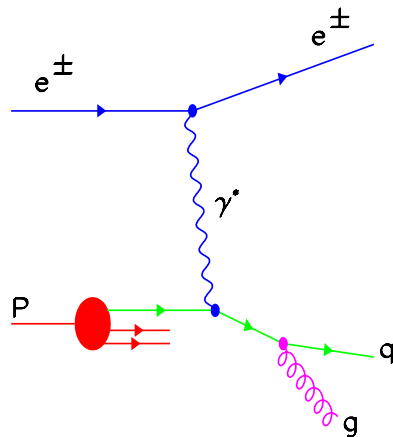
$$y = 1 - \frac{E'_e}{2E_e} (1 - \cos \vartheta)$$

$$Q^2 = 2E'_e E_e (1 + \cos \vartheta)$$

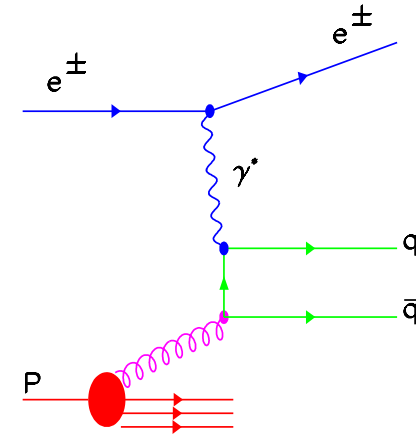
$$x = \frac{Q^2}{ys}$$

where  $\sqrt{s}$  is the eP centre of mass energy

# QCD Improved QPM



(a)



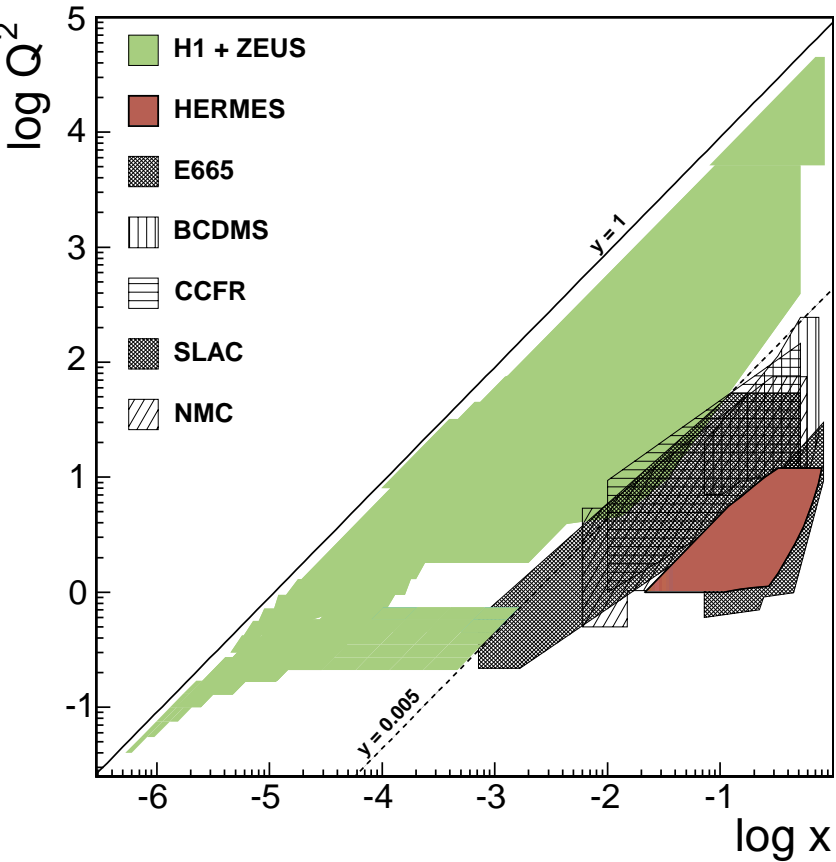
(b)

Leading order  $\mathcal{O}(\alpha_s)$  modifies QPM picture:

There are 2 contributions:

- QCD Compton - the quark radiates a gluon before or after being struck by the virtual photon
- Boson Gluon Fusion - the virtual photon & a gluon inside the proton produce a quark-antiquark pair

# HERA kinematic range



HERA extends the kinematic reach of previous DIS expt:

- $Q^2$  in range  $10^{-1}$  to  $10^5 \text{ GeV}^2$
- $x$  down to  $10^{-6}$

extension by two orders of magnitude in both  $x$  and  $Q^2$

## DIS NC X-section

$$\frac{d^2 \sigma^{e^\pm p}}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} (Y_+ F_2(x, Q^2) - y^2 F_L(x, Q^2) - Y_- x F_3(x, Q^2))$$

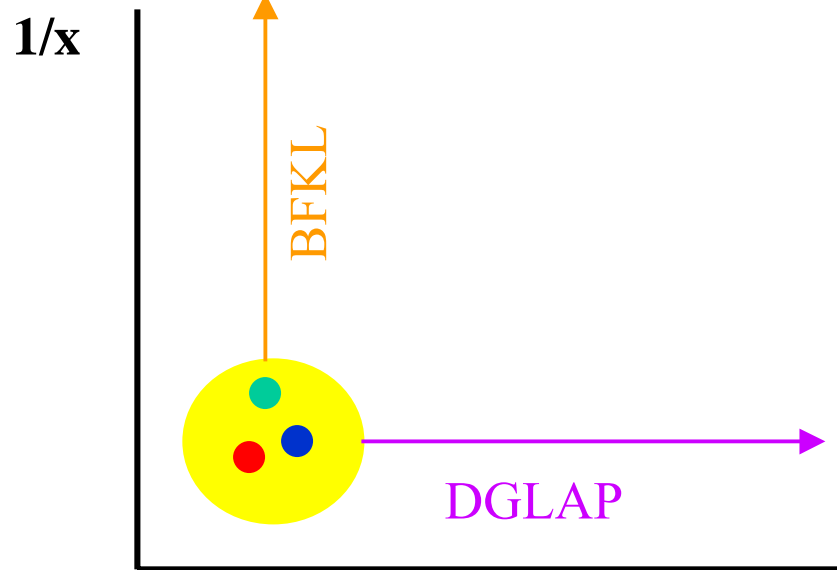
where  $Y_\pm = 1 \pm (1-y)^2$

$$F_2 = e_i^2 (q(x) + \bar{q}(x)) \quad (\text{in QPM})$$

$F_L$  – long. str. fnc; important only for  $y > 0.6$

$F_3$  – arises from Z-exchange;

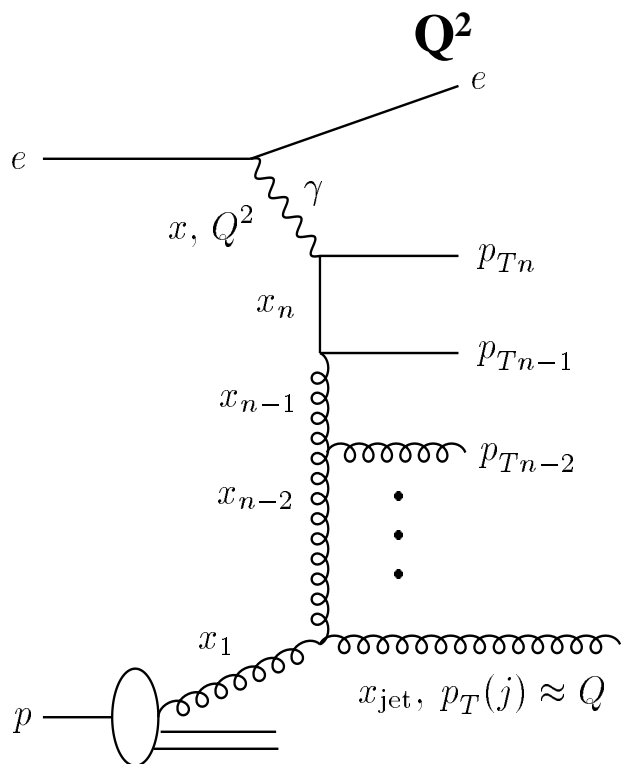
negligible for  $Q^2 < 5000 \text{ GeV}^2$



# DGLAP vs BFKL

## Parton Evolution in DIS

DGLAP	BFKL
Strong $p_T$ ordering	No $p_T$ ordering
Soft frac. mom. ordering	Strong frac. mom. ordering



# NLO QCD Fits

DGLAP predicts  $Q^2$  evolution of  $F_2(x, Q^2)$  for given parton densities  
at  $Q^2 = Q_0^2$

## H1 fit

$$Q_0^2 = 1 \text{ GeV}^2$$

$$\text{gluon : } xg(x, Q_0^2) = A_g x^{B_g} (1-x)^{C_g}$$

$$\text{valence } u_v : xu_v(x, Q_0^2) = A_u x^{B_u} (1-x)^{C_u} (1 + D_u x^{E_u})$$

$$\text{valence } d_v : xd_v(x, Q_0^2) = A_d x^{B_d} (1-x)^{C_d} (1 + D_d x^{E_d})$$

$$\text{sea : } xS(x, Q_0^2) = A_s x^{B_s} (1-x)^{C_s}$$

$$\text{strange quarks: } \bar{s} = \bar{u}/2$$

assume  $\bar{u} - \bar{d}$  param. from MRS

$$\alpha_s(M_Z) = 0.118$$

## ZEUS fit

$$Q_0^2 = 7 \text{ GeV}^2$$

$$\text{gluon : } xg(x, Q_0^2) = A_g x^{\delta_g} (1-x)^{\eta_g} (1 + \gamma_g x)$$

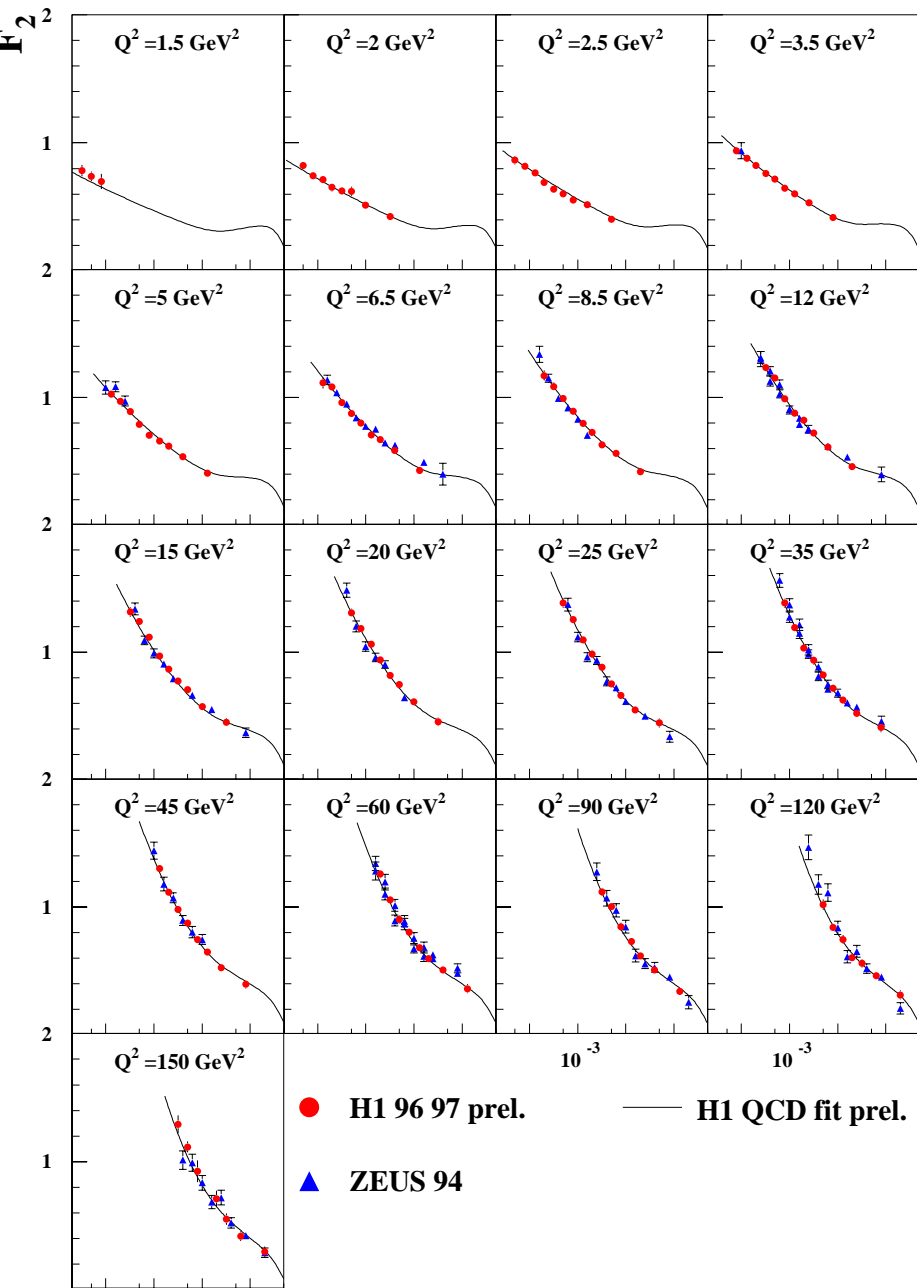
$$\text{sea : } xS(x, Q_0^2) = A_s x^{\delta_s} (1-x)^{\eta_s} (1 + \gamma_s x + \epsilon_s \sqrt{x})$$

$$u-d \text{ difference : } x\Delta_{ud}(x, Q_0^2) = A_{\Delta}^{\delta_{\Delta}} (1-x)^{\eta_{\Delta}}$$

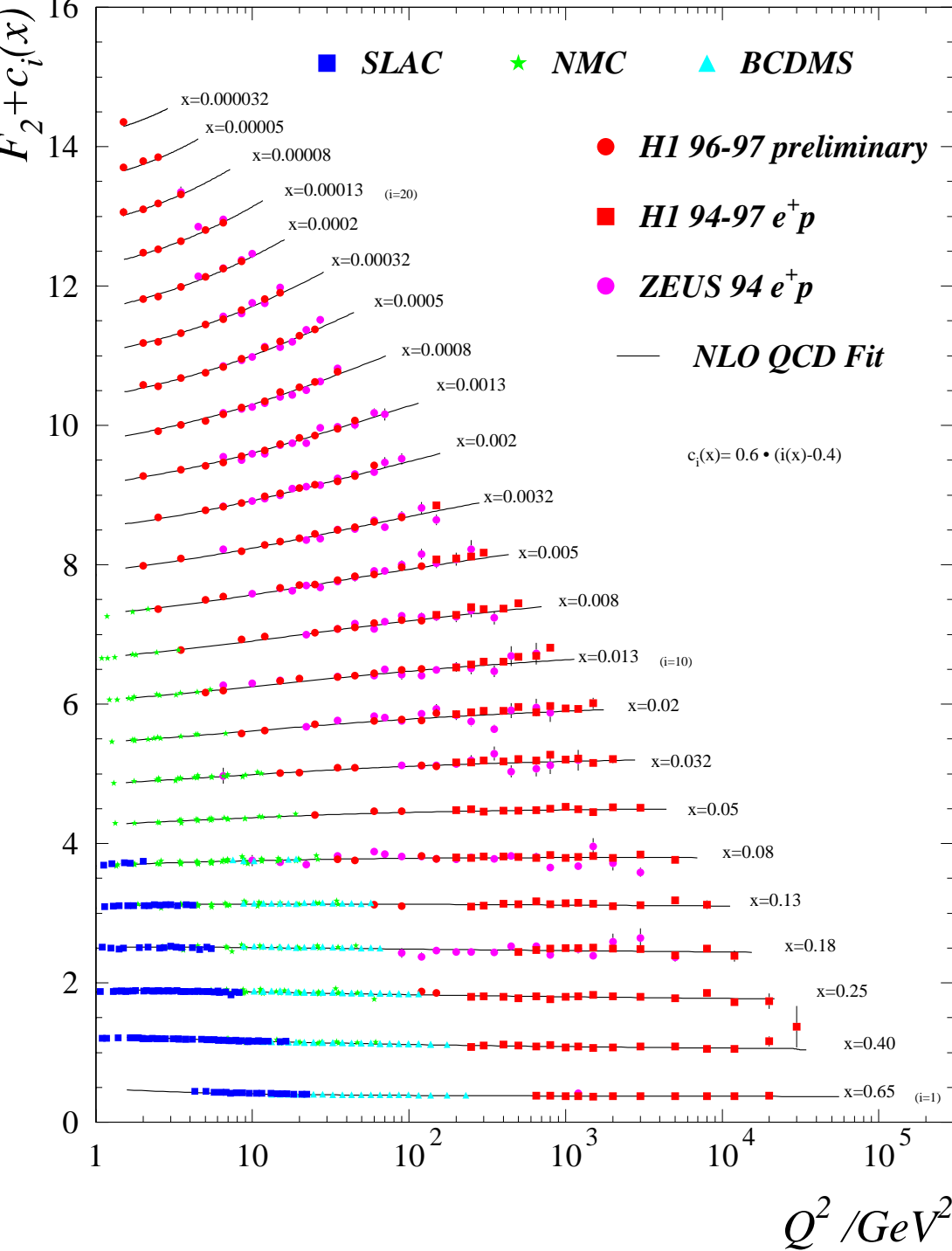
strange quark assumed 20% of sea  
valence quarks from MRS(R2)

$$\alpha_s(M_Z) = 0.118$$

- fixed flavour scheme - 3 light flavours, heavy flavours in NLO via BGF
- momentum sum rule
- quark counting rules



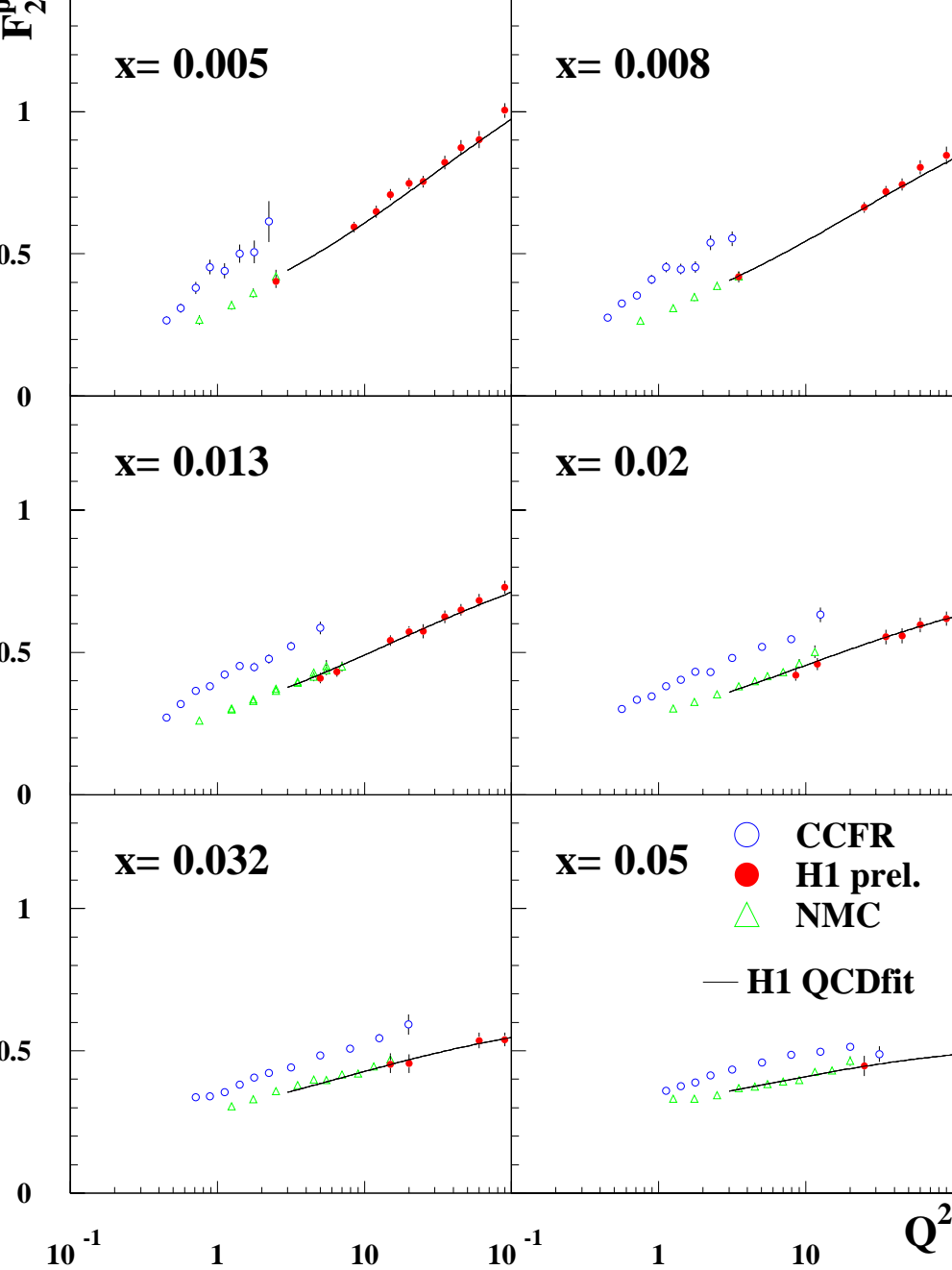
- strong rise of  $F_2$  at low  $x$
- good agreement between expts
- systematically dominated (2-3%) up to  $Q^2 \approx 1000 \text{ GeV}^2$



NLO DGLAP fit gives good description of the HERA & fixed target data

Scaling violation well interpreted by QCD

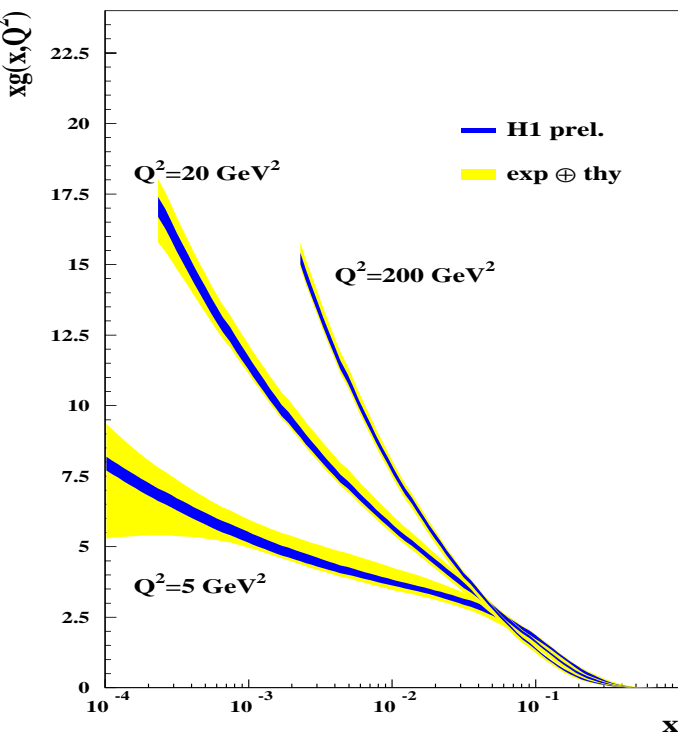
No indication of  $(\log 1/x)^n$  corrections in HERA regime



Long standing controversy between  $\mu p$  (NMC, E665) and  $\nu N$  (CCFR) data

H1 data overlap and extrapolate well to  $\mu p$  data

CCFR data being re-analysed, with new treatment of charm and shadowing



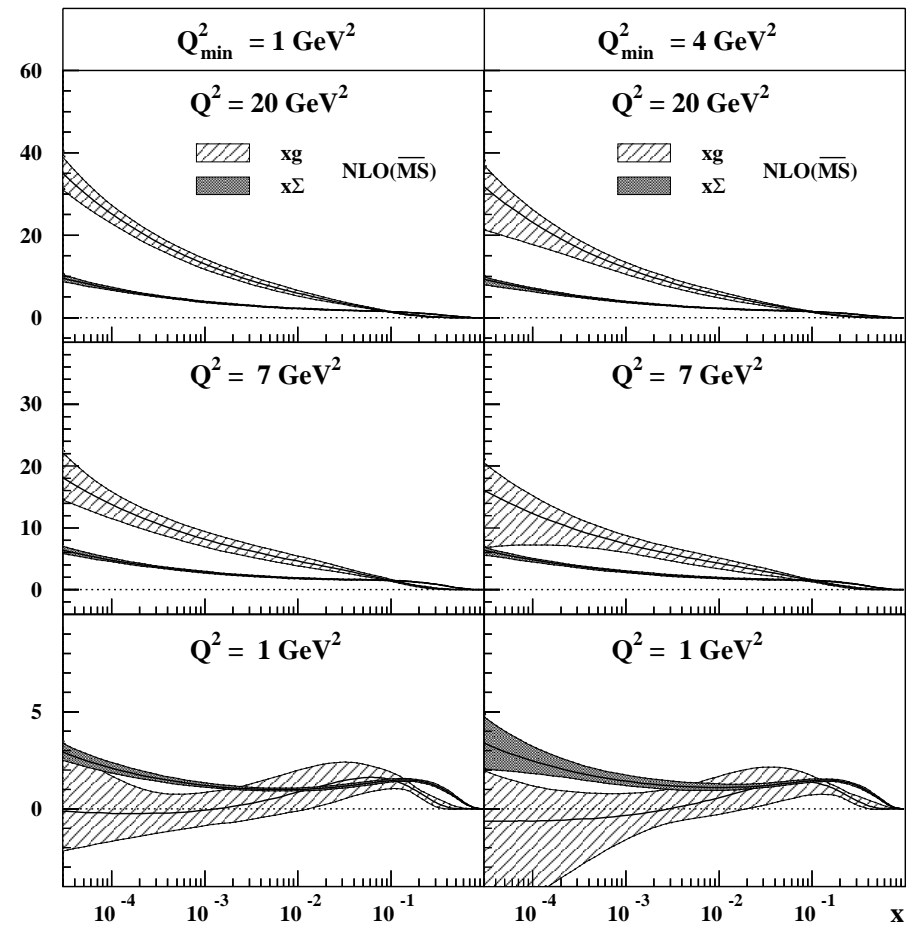
Strong rise in  $xg$  as  $x \downarrow$

10% uncertainty at  $Q^2=20$   $\text{GeV}^2$  &  $x=5 \times 10^{-5}$

$F_2$  rise at low  $x$  not completely driven by the increase of gluon density from parton splitting

# Gluon Determination

## ZEUS 1995



# Extraction of $F_L$

remember  $Y_{\pm} = 1 \pm (1 - y)^2$

- subtraction method

$$\sigma_r = \left( \frac{xQ^4}{2\pi\alpha^2 Y_+} \right) \frac{d^2\sigma}{dx dQ^2} = F_2 - \frac{y^2}{Y_+} F_L$$

measure  $\sigma_r$  at high  $y$  - for  $Q^2 \ll M_Z^2$

fitted  $F_2$  at low  $y$  extrapolated to high  $y$  & subtract out  $F_L$

- ★ derivative method

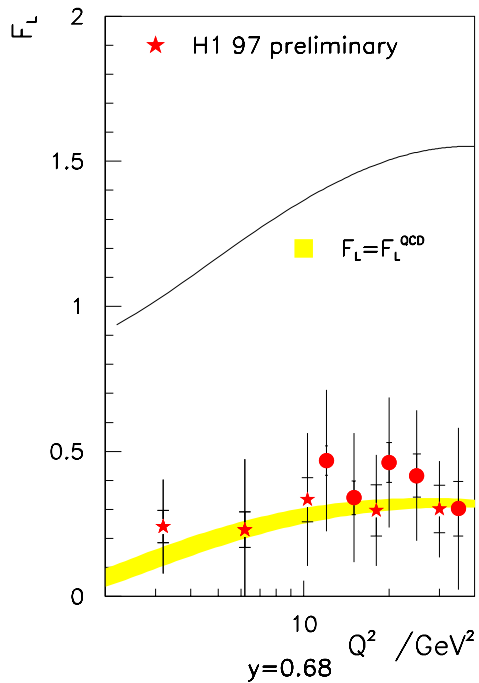
$$\frac{d\sigma_r}{d \log y} = -\frac{dF_2}{d \log y} - 2y^2 \frac{2-y}{Y_+^2} F_L + \frac{y^2}{Y_+} \frac{dF_L}{d \log y}$$

assume  $\frac{dF_2}{d \log y} = A \log y + B$

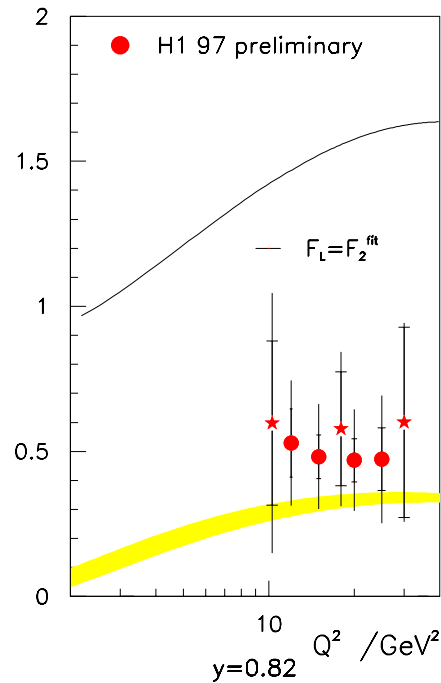
straight line fit to  $d\sigma_r/d \log y$   $Q^2$  bins at  $y < 0.2$

# Extraction of $F_L$

★ derivative



● subtraction

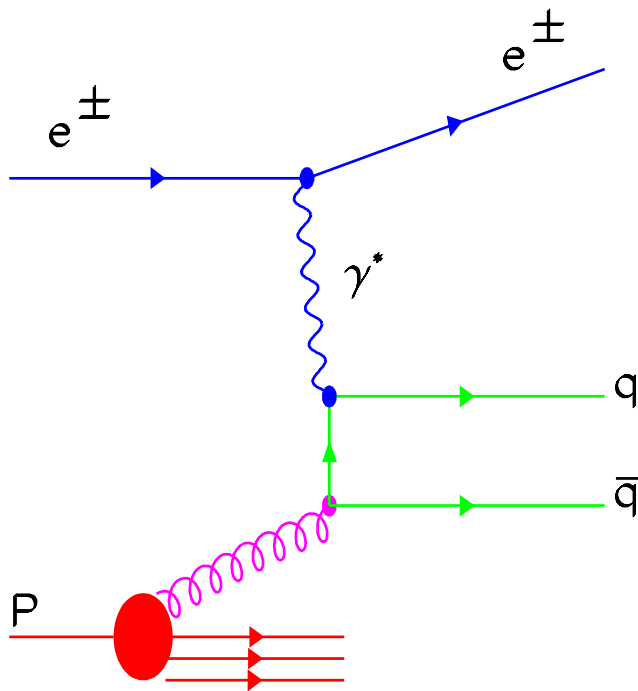


$F_L$  compatible with QCD predictions

Both methods in agreement

Direct measurement needs different beam energies or ISR events

# Investigation of gluon via charm production



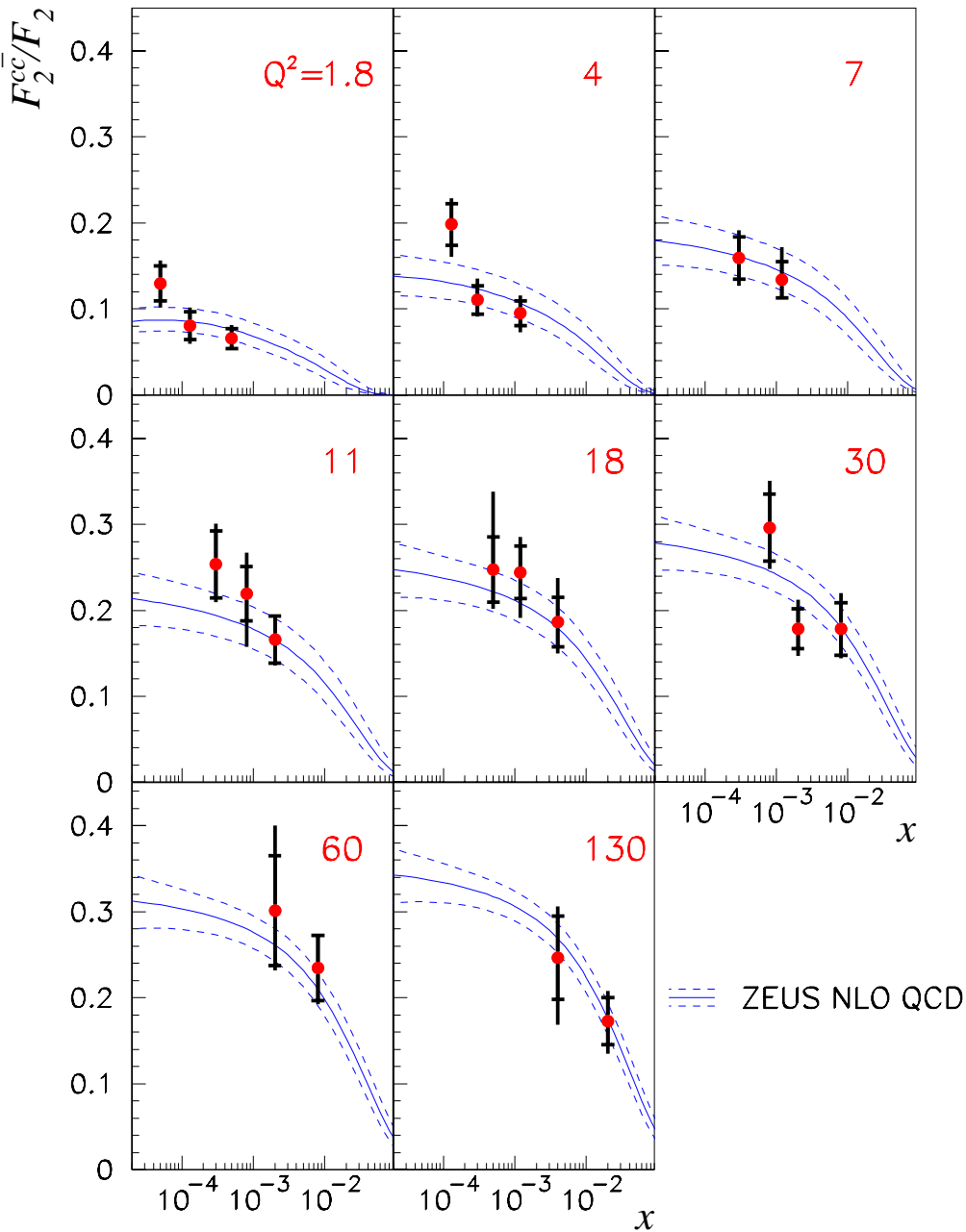
(b)

- Charm production dominated by BGF diagram

probe of the gluon density

- investigate via  $D^*$  production

# ZEUS 1996-97



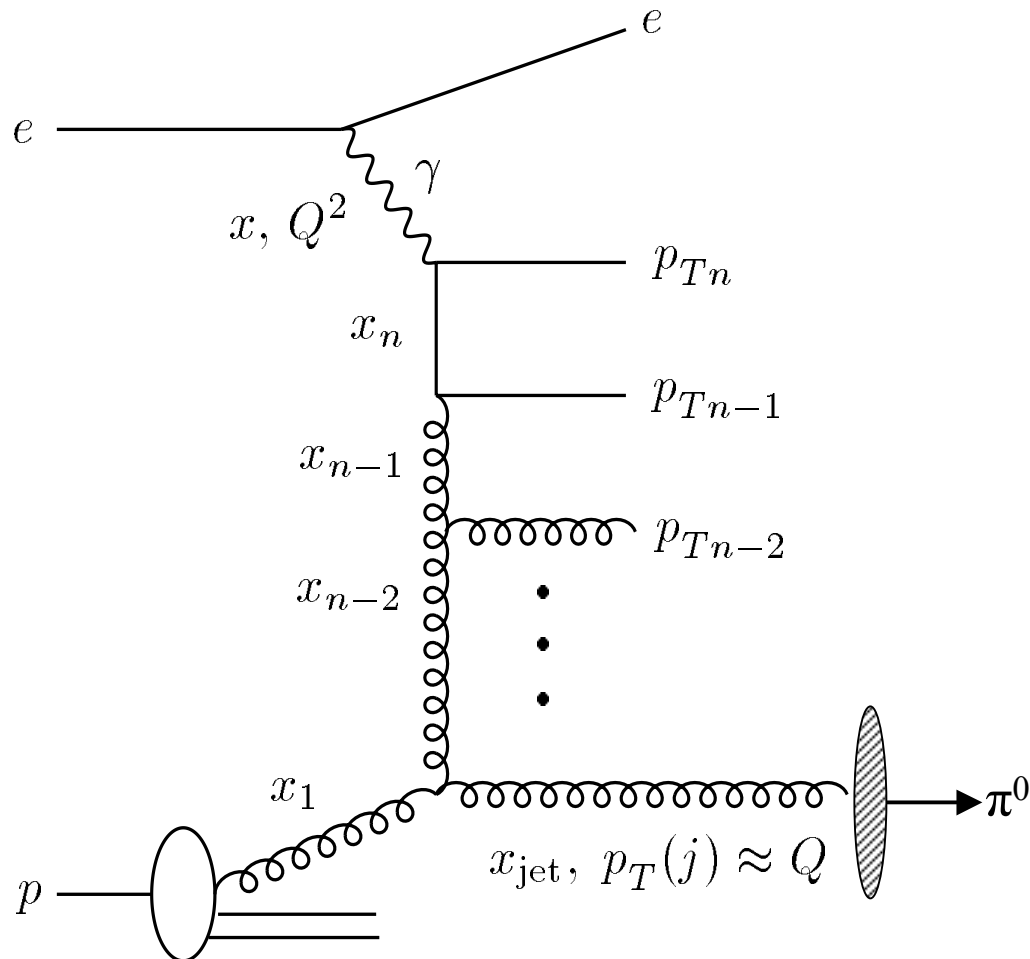
$F_c \approx 25\%$  at  
low  $x$  & high  $Q^2$

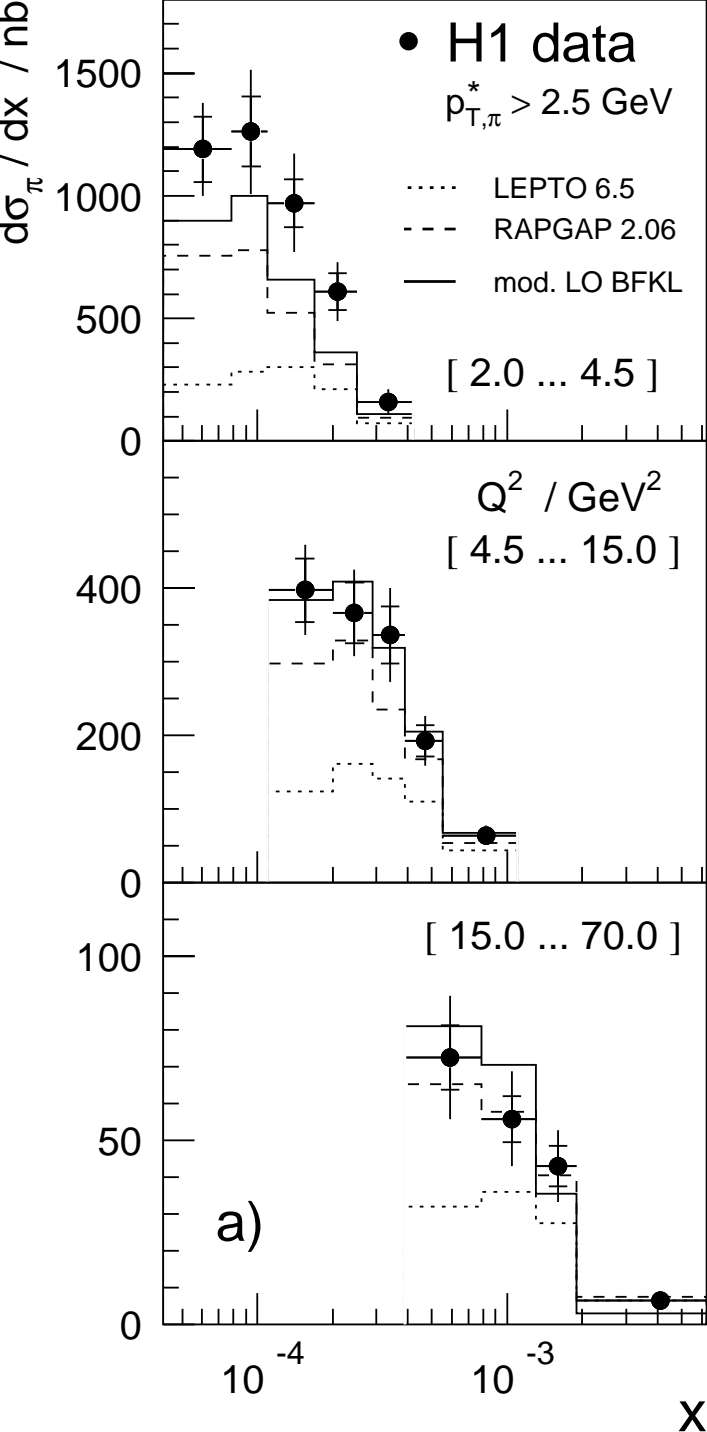
$F_c$  at  $Q^2=1.8 \text{ GeV}^2$  is  
 $\approx 10\%$

$F_c$  rising quicker than  $F_2$

# DGLAP vs BFKL – forward $\pi^0$ production

BFKL predicts greater forward  $\pi^0$  production than DGLAP expectation





MC models implementing DGLAP evolution fail to describe  $\pi^0$  data

adding resolved component to incoming virtual photon improves description of data

BFKL formalism gives good description of data. Still question marks over absolute normalisation

# Event Shape Variables

Any ‘infra-red’ safe event variable  $\langle F \rangle$  can be written as

$$\langle F \rangle = \langle F \rangle^{\text{pert}} + \langle F \rangle^{\text{pow}}$$



Log change of the strong  
coupling const  $\propto 1/\log(Q)$

Power corrections or  
hadronisation effects  $\propto 1/Q$

$1/Q$  corrections not necessarily related to hadronisation

**BUT**

soft gluon phenomenon at small momentum scales

$$\langle F \rangle^{\text{pert}} = c_1(x, Q) \alpha_s(\mu_R) + \left[ c_2(x, Q) + \frac{\beta_0}{2\pi} \log\left(\frac{\mu_R}{Q}\right) c_1(x, Q) \right] \alpha_s^2(\mu_R)$$

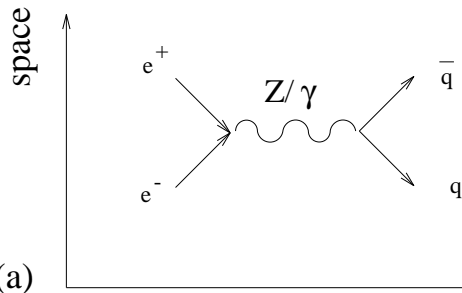
$c_1$  &  $c_2$  are coefficients in the  $\overline{\text{MS}}$  scheme

$\mu_R$  is the renormalisation scale taken to be  $Q$

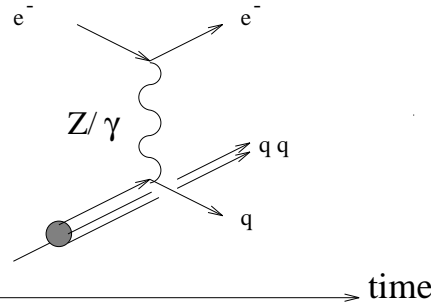
# The Breit Frame

## 'Brickwall' frame

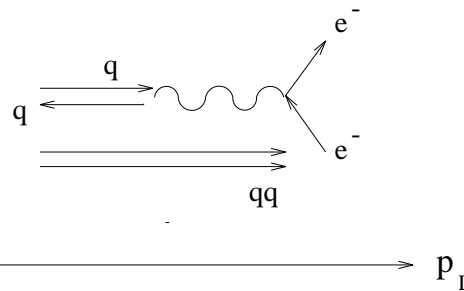
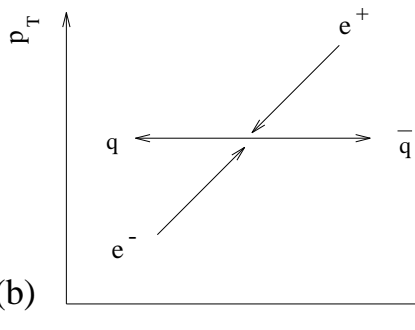
Electron-positron Annihilation



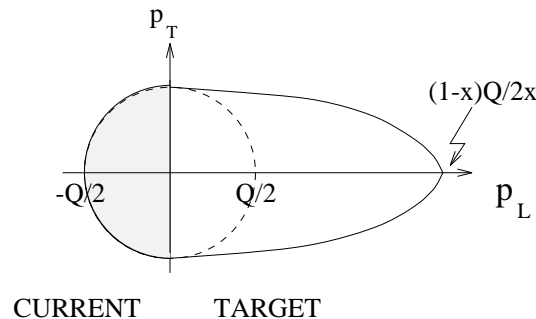
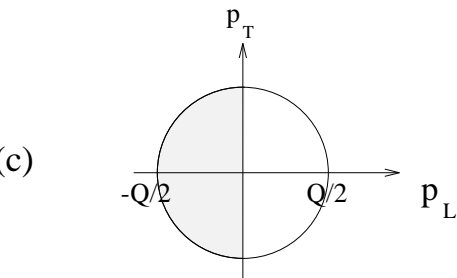
DIS in the Breit Frame



Phase space for  $e^+e^-$  annihilation evolves with  $Q/2 = \sqrt{s}/2$



Current hemisphere of Breit frame evolves as  $Q/2$



Current region  $\equiv e^+e^-$  annihilation

Thrust  $T_C$  or  $T_m$

$$\tau_C = 1 - T_C = 1 - \max_h \frac{|p_h \cdot n_T|}{|p_h|}; n_T \equiv \text{thrust axis}$$

Thrust  $T$  or  $T_Z$

$$\tau = 1 - T = 1 - \frac{|p_h \cdot n|}{|p_h|}; n \equiv \text{hemisphere axis}$$

Jet mass  $\rho$

$$\rho = \frac{M^2}{(2E_{\text{tot}})^2} = \frac{\left( \sum_h p_h \right)^2}{4 \left( \sum_h E_h \right)^2}$$

C parameter

$$C = 3(\lambda_1 \lambda_2 + \lambda_2 \lambda_3 + \lambda_3 \lambda_1)$$

where  $\lambda_i$  are the eigenvalues of the momentum tensor

$$\Theta_{jk} = \frac{\sum_h p_{j_h} p_{k_h}}{\sum_h |p_h|^2}$$

$y_{fj}$  &  $y_{k_t}$  are transition values for  $(2+1) \rightarrow (1+1)$  jets for the factorisable JADE algo. & the  $k_t$  algo. respectively

Jet Broadening  $B$

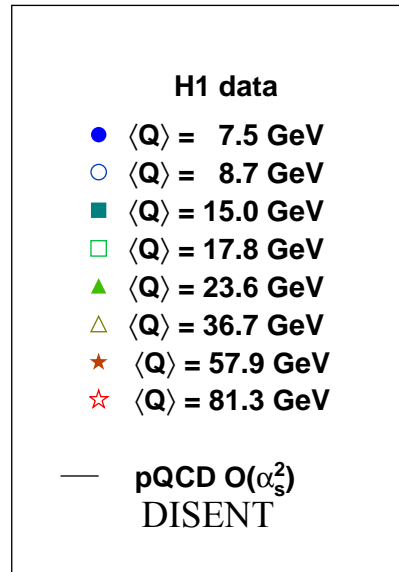
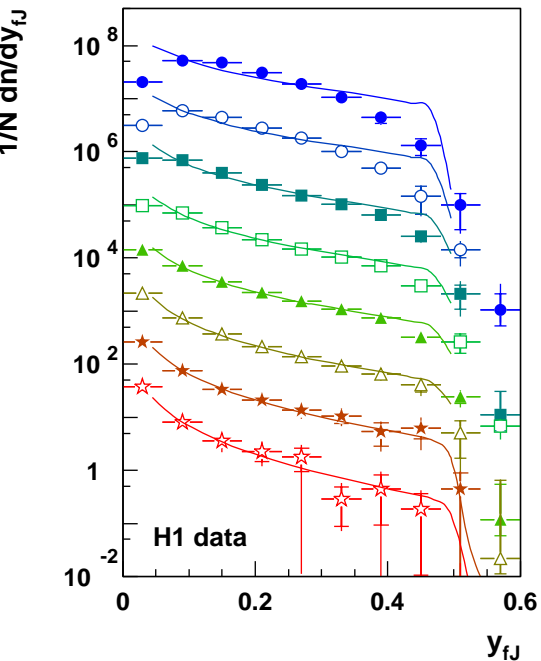
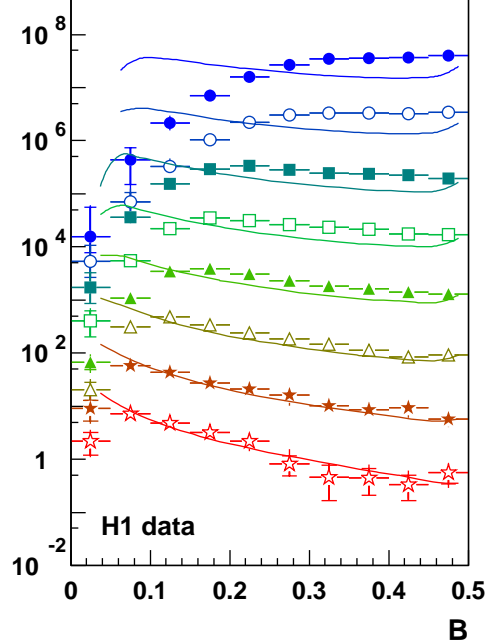
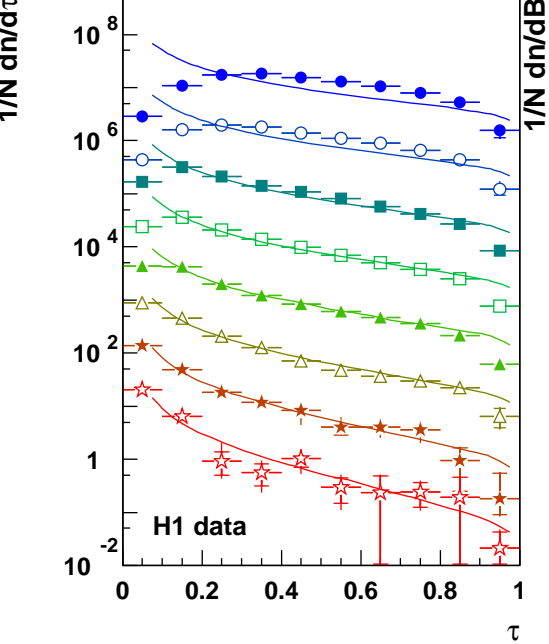
$$B = \frac{\sum_h |p_h \times n|}{2 \sum_h |p_h|} = \frac{\sum_h |p_{\perp h}|}{2 \sum_h |p_h|}$$

# Data vs NLO

Events more pencil like as  $Q^2 \uparrow$

Non-pert corrections decrease as  $Q^2 \uparrow$

Non-pert correction for jet variables smaller over all  $Q^2$



$$\langle F \rangle^{\text{pow}} = a_F \frac{32}{3\pi^2} \frac{M}{p} \left( \frac{\mu_I}{Q} \right)^p \left[ \bar{\alpha}_{p-1}(\mu_I) - \alpha_s(Q) - \frac{\beta_0}{2\pi} \left( \ln \frac{Q}{\mu_I} + \frac{K}{\beta_0} + \frac{1}{p} \right) \alpha_s^2(Q) \right]$$

$\beta_0, K$  are constant dependent on number of flavours

$\mu_I$  - 'infra-red' matching scale,  $\mu_I = 2 \text{ GeV}$

$a_F, p$  - calculable coeff dependent on observable  $F$

$p = 1$  except for  $y_{k_i}$  where  $p = 2$

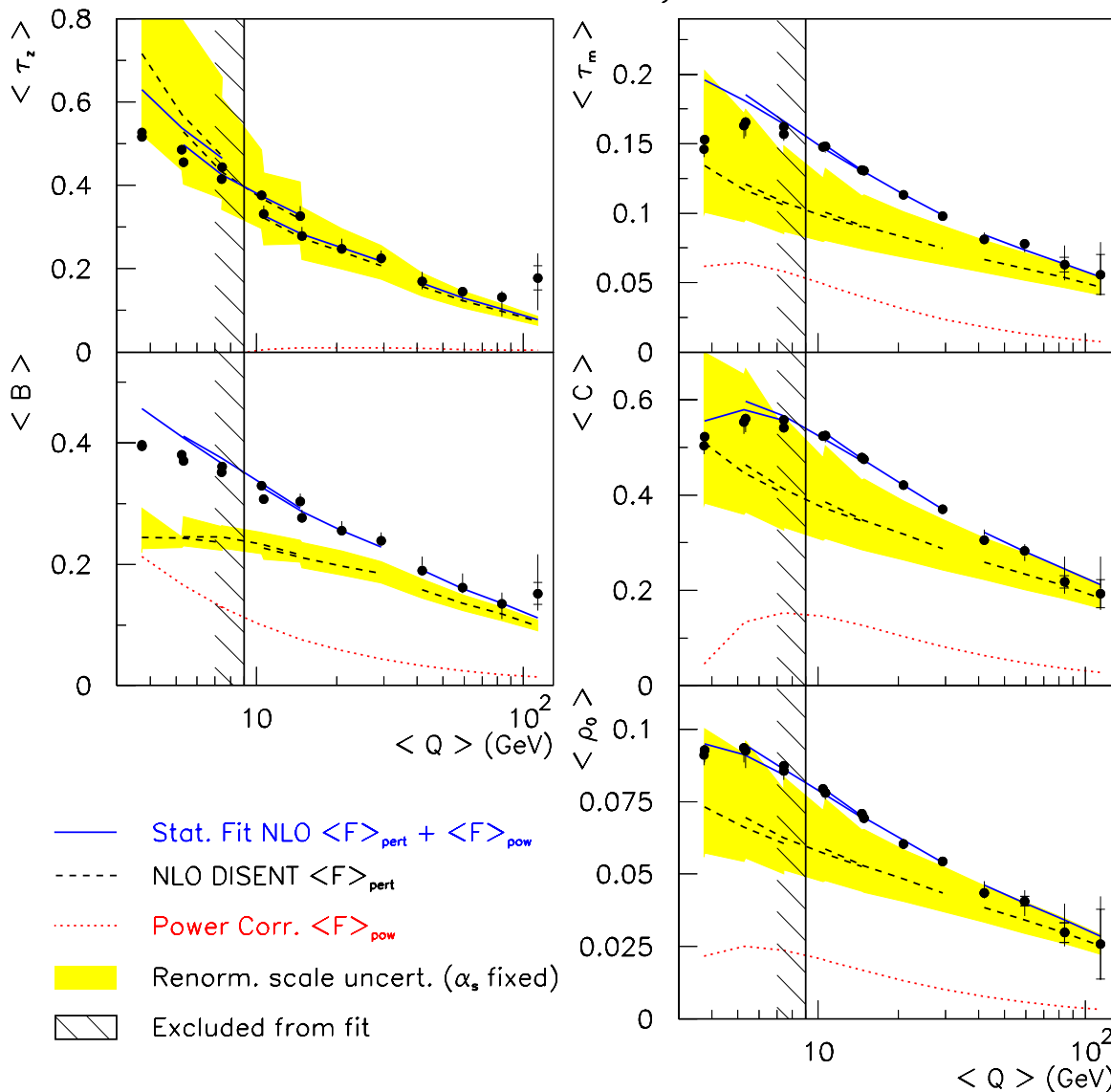
For  $B, a_F = F(\alpha_{s-CMW}(Qe^{-3/4}), N_f)$

$M \approx 1.43$  - 2-loop correction (Milan factor)

$\bar{\alpha}_{p-1}$  - an universal, non-pert. effective strong coupling below  $\mu_I$

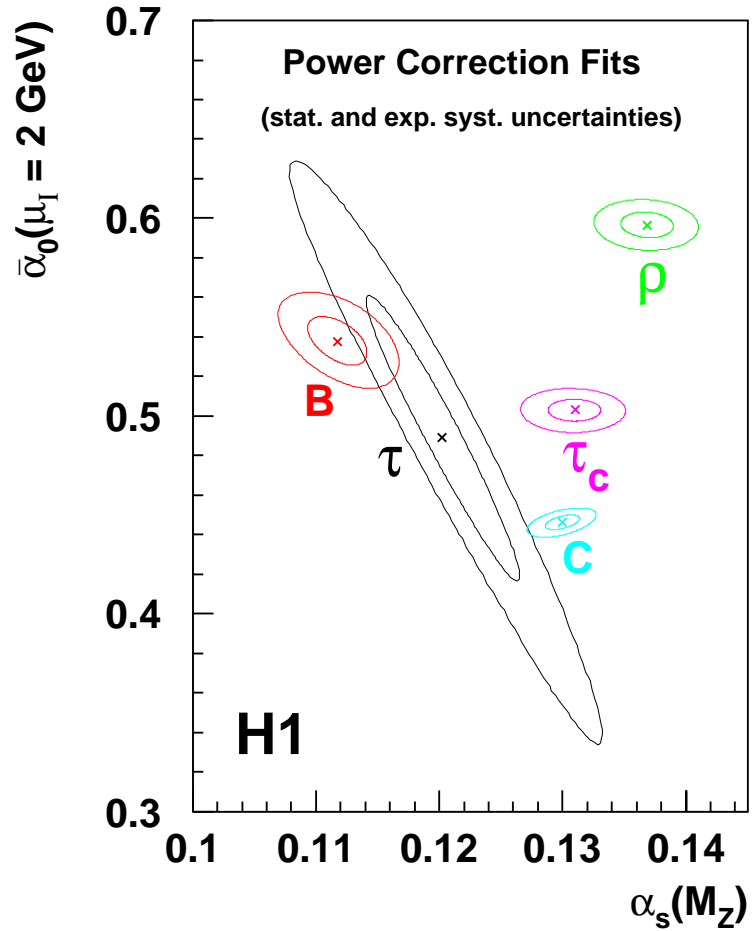
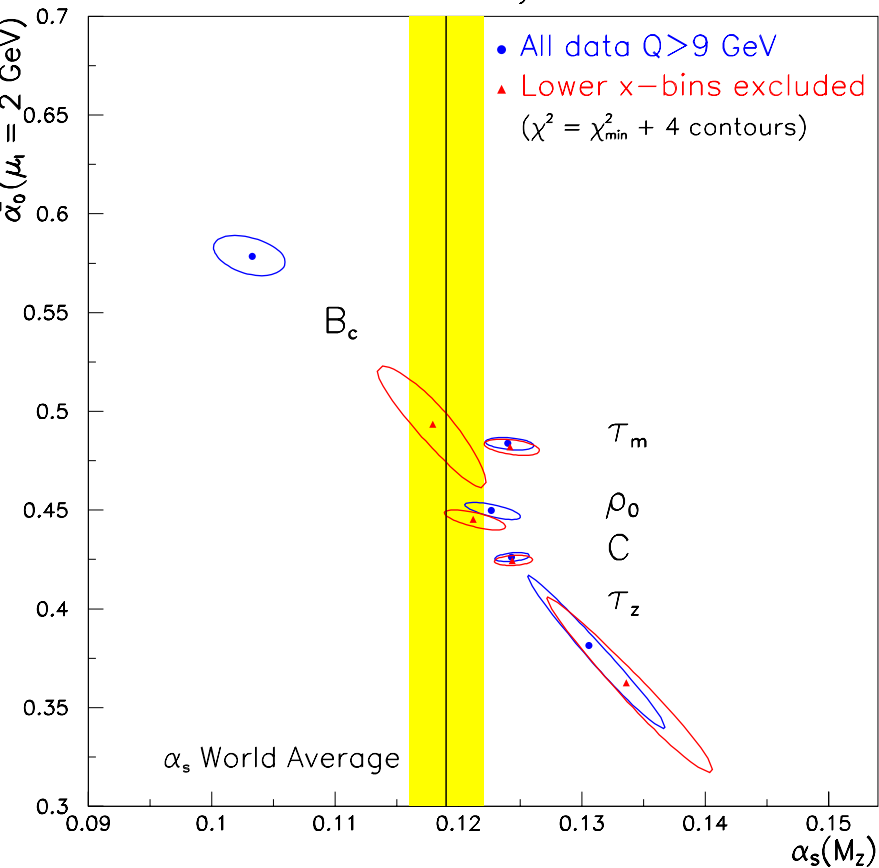
Power correction  
fit

ZEUS Preliminary 1995–97



Reasonable fit to data

Closer examination shows NLO calculation (for B in particular) has wrong x dependence

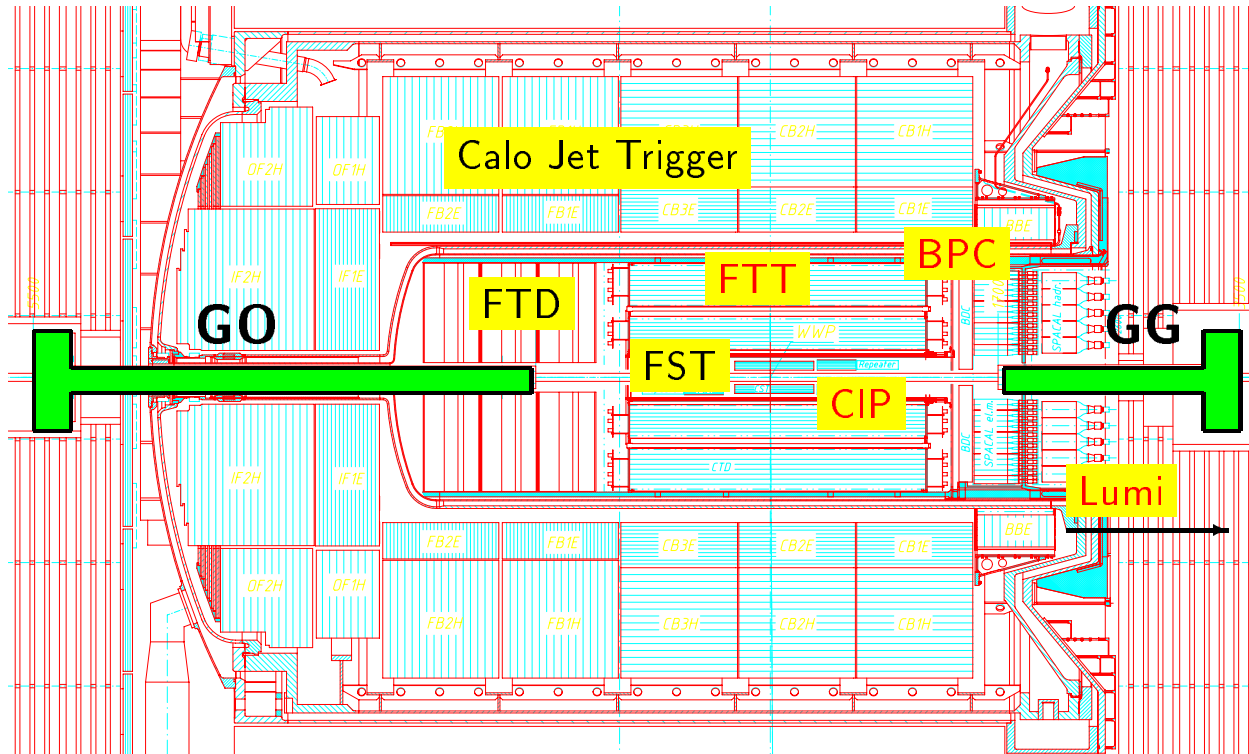


x-dependence on result, particular for  $B$  and  $T$

H1 & ZEUS consistent (exception  $\rho$ , where different def<sup>n</sup> used)

$y_{fJ}$  power correction coeff  $a_F$  not compatible with data

# Upgrade of HERA & Detectors



- Magnets around beamline (including inside detector vol)

- $\beta$  functions reduced by factors of 4-5

- increased currents

- factor of 5 increase in luminosity

- e-polarisation

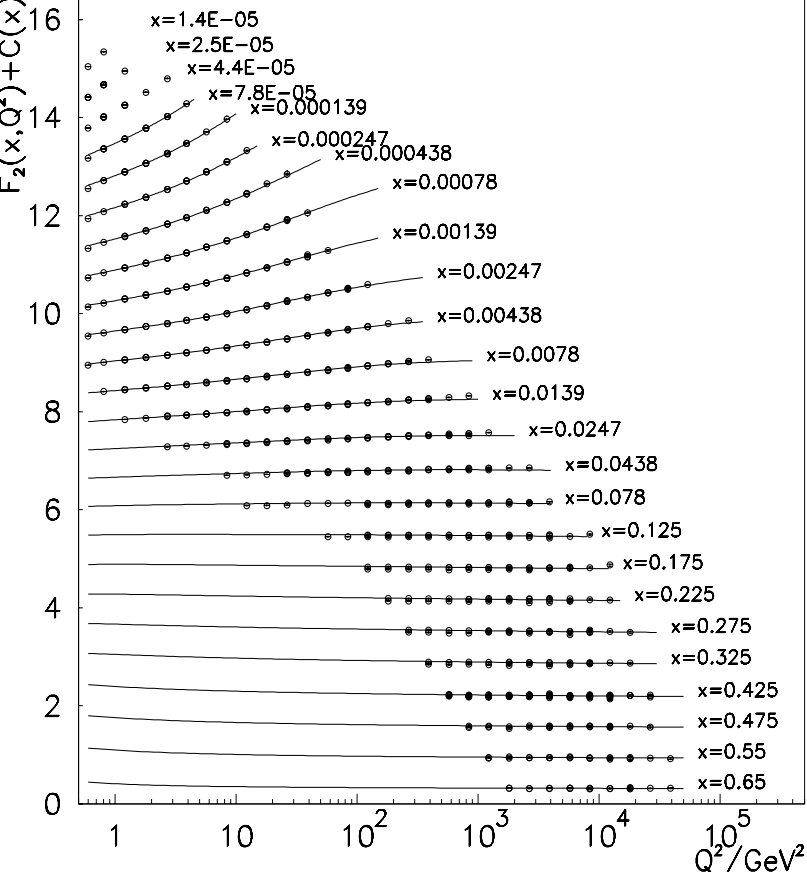
- forward tracking upgrade

- new lumi detectors

- microvertex detector(ZEUS)

- improved tracking trigger (H1)

- ⋮



simulation of  $F_2$  with  $1 \text{ fb}^{-1}$  at HERA

measurements of  $F_c$  and  $F_b$  to 5% & 10% respectively

stringent tests of QCD evolution

important expt. input to future hadron colliders !

precision on  $\alpha_s(M_Z)$  0.001

gluon density extraction to 1%

# Summary

- structure func precision a few %
- DGLAP evolution OK down to  $Q^2 \approx 1 \text{ GeV}^2$
- $F_c$  up to 25% of  $F_2$
- an indirect measurement of  $F_L$
- event shapes in reasonable agreement with NLO & power corrections. Still outstanding questions
- HERA high luminosity running deliver  $1 \text{ fb}^{-1}$  per expt. during 2001  $\rightarrow$  2006