Physics at LHC
Cracow, Poland
July $3^{\rm rd} - 8^{\rm th}$, 2006

Hadronic final states and QCD studies in ep collisions

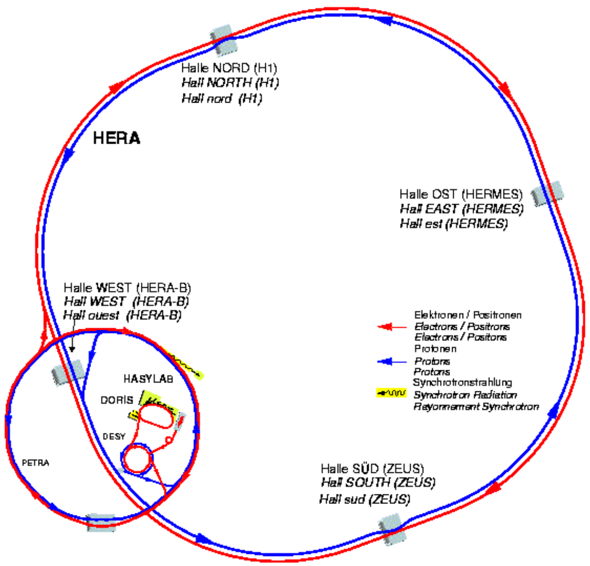
Claudia Glasman
Universidad Autónoma de Madrid

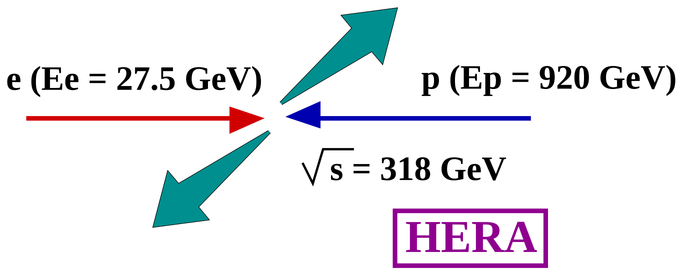
from

ZEUS Collab.

H1 Collab.

at

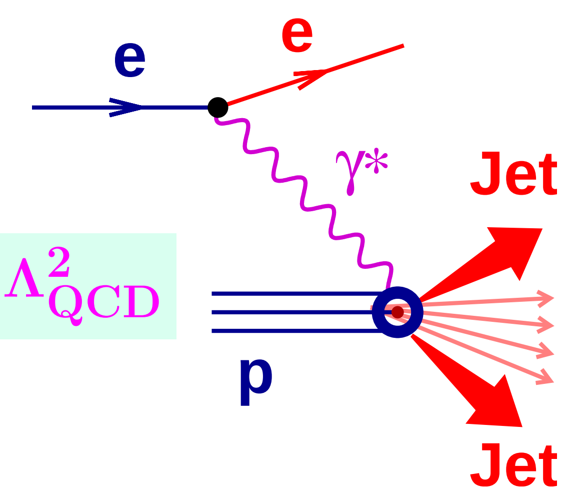


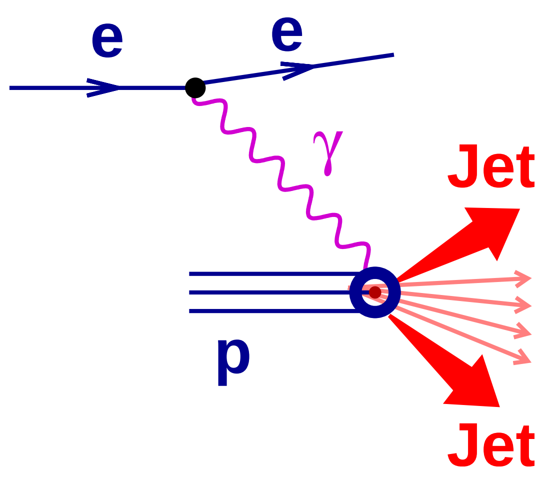


Introduction

- Jet production in pp collisions will be copious at LHC
- Main background to searches in hadronic channels: QCD multijet production
- Also, luminosity of colliding particles (partons from protons) governed by QCD (proton PDFs)
- ep collider HERA: very suitable enviroment to do precision studies of QCD
  $\longrightarrow$ tests of QCD in hadronic-induced reactions (as opposed to e^+e^- at LEP)
  $\longrightarrow$ but cleaner than $p\bar{p}$ at TeVatron
- The main sources of jets at HERA are:

$Q^2 \gg \Lambda^2_{\rm QCD}$





$Q^2 \approx 0$

NC deep inelastic scattering (DIS)

$$ep \longrightarrow e + \mathrm{Jet}\ (+\mathrm{Jet}) + X$$

photoproduction (γp)

$$ep \longrightarrow e + \mathrm{Jet}\ (+\mathrm{Jet}) + X$$

How do we study QCD at HERA ?

- **We measure:**

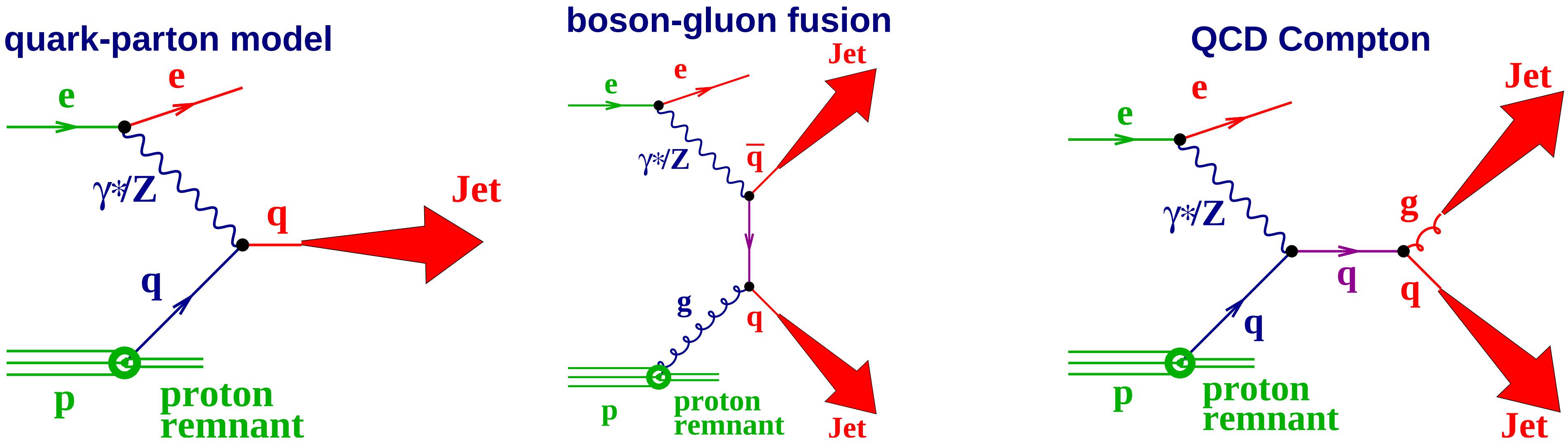
structure functions	heavy quark production	jet production	particle production
↓	↓	↓	↓
F_2^p, F_2^{CC}, F_2^D, F_2^c, ...	charm beauty	jet cross sections jet substructure forward jets	hadrons photons event shapes

to obtain results on:

- — proton, photon, pomeron structure: parton densities
- — tests of pQCD: up to which extent does pQCD describe jet dynamics?
- — parton dynamics at low x: breakdown of DGLAP?
- — measurements of α_s: the fundamental parameter of the theory
- — color dynamics: color factors, subprocesses
- — non-perturbative effects: can they be described from first principles?

Jet production in neutral current deep inelastic ep scattering

- Jet production in neutral current deep inelastic ep scattering up to $\mathcal{O}(\alpha_s)$:



Kinematics:

– momentum transfer:

$$Q^2 = -q^2 = -(k - k')^2$$

– Bjorken x: $x = \frac{Q^2}{2P\cdot q}$

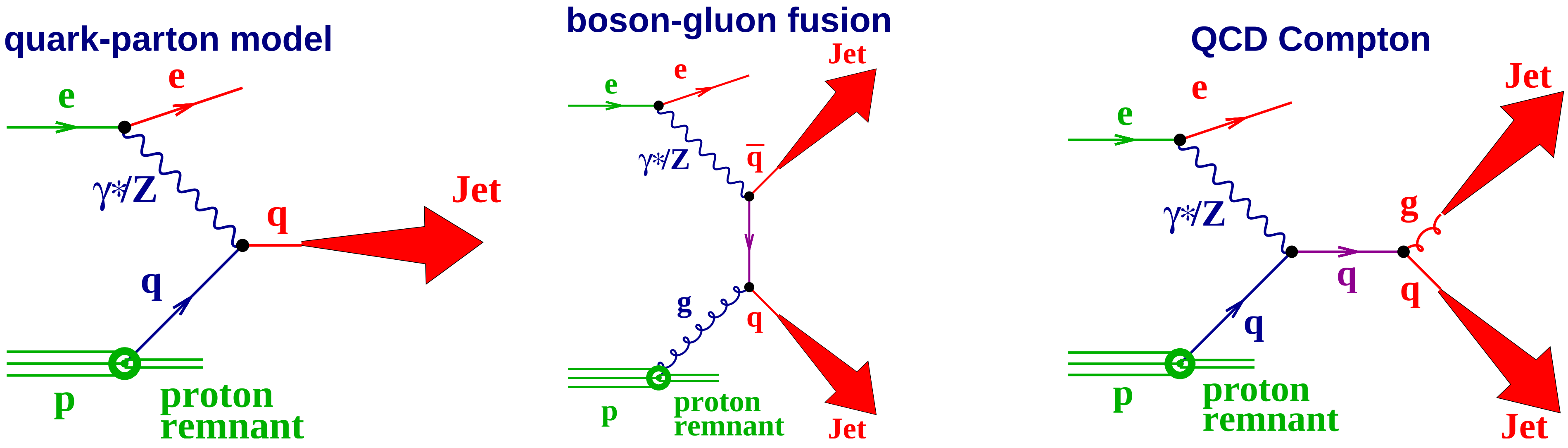
(parton momentum fraction for QPM)

– inelasticity:

$$y = \frac{P\cdot q}{P\cdot k} = 1 - \frac{E'_e(1-\cos\theta_e)}{2E_e}$$

Jet production in neutral current deep inelastic ep scattering

- Jet production in neutral current deep inelastic ep scattering up to $\mathcal{O}(\alpha_s)$:



- Jet production cross section:

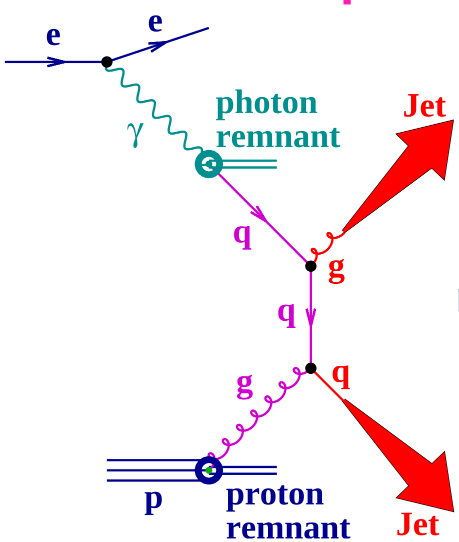
$$d\sigma_{\text{jet}} = \sum_{a=q,\bar{q},g} \int dx\, f_a(x, \mu_F)\, d\hat{\sigma}_a(x, \alpha_s(\mu_R), \mu_R, \mu_F)$$

– f_a: parton a density in the proton, determined from experiment
  $\longrightarrow$ long-distance structure of the target

– $\hat{\sigma}_a$: subprocess cross section, calculable in pQCD
  $\longrightarrow$ short-distance structure of the interaction

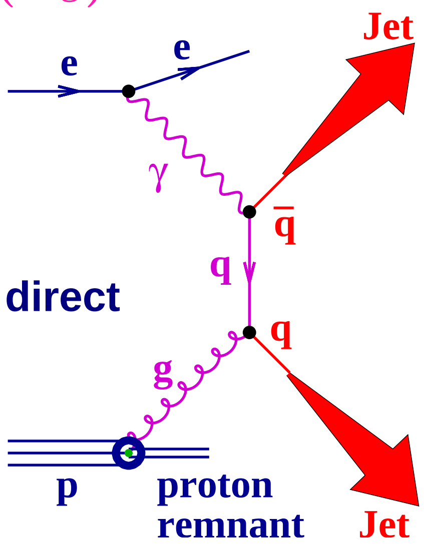
Jet production in photoproduction

- **Jet production in photoproduction up to $\mathcal{O}(\alpha_s)$:**



resolved

direct



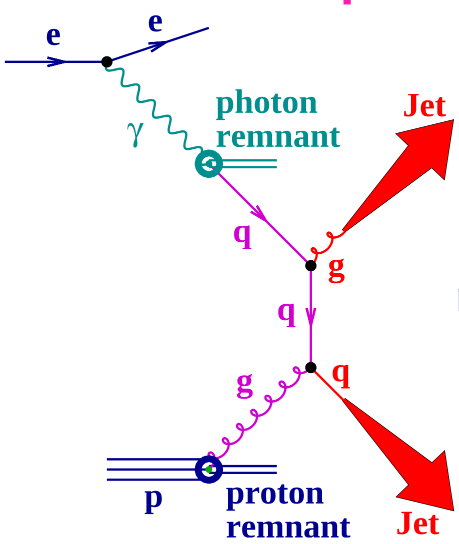
Kinematics:

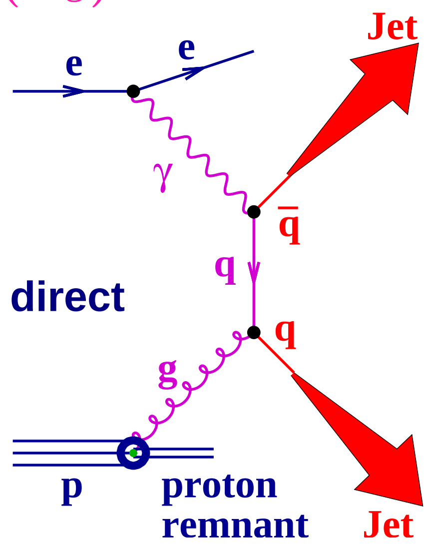
– $Q^2 \approx 0$

– **total hadronic cms energy:** $W^2 = ys$

Jet production in photoproduction

- **Jet production in photoproduction up to $\mathcal{O}(\alpha_s)$:**





- **Jet production cross section:**

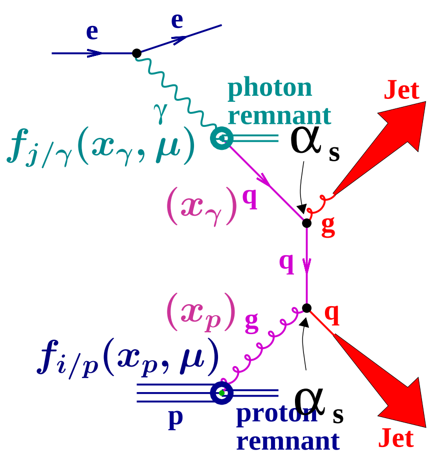
$$d\sigma_{\text{jet}} = \sum_{i,j} \int_0^1 dy\, dx_\gamma\, dx_p\, f_{\gamma/e}(y)\ \ f_{i/\gamma}(x_\gamma, \mu_{F_\gamma})\ \ f_{j/p}(x_p, \mu_{F_p})\ d\hat{\sigma}_{i(\gamma)j}(i(\gamma)j \rightarrow \text{jet jet})$$

– $f_{j/p}(f_{i/\gamma})$: parton density in the proton (photon)
$\longrightarrow$ long-distance structure of the target

– $\hat{\sigma}_{i(\gamma)j}$: subprocess cross section, calculable in pQCD
$\longrightarrow$ short-distance structure of the interaction

Jet search in ep collisions

- Jet search in e^+e^- annihilations is simple:
  - $\longrightarrow$ initial state: only leptons; final state: arising uniquely from short-distance interaction (all hadrons in final state associated with hard process)
  - $\longrightarrow$ best frame: centre-of-mass system = LAB
  - $\longrightarrow$ variables invariant under rotations: energies and angles
  - $\longrightarrow$ distance between hadrons: angular separation
- Jets in hadronic collisions are not as easily identified because jets carry only a fraction of the available energy and are accompanied by several soft hadrons not correlated with the hard interaction

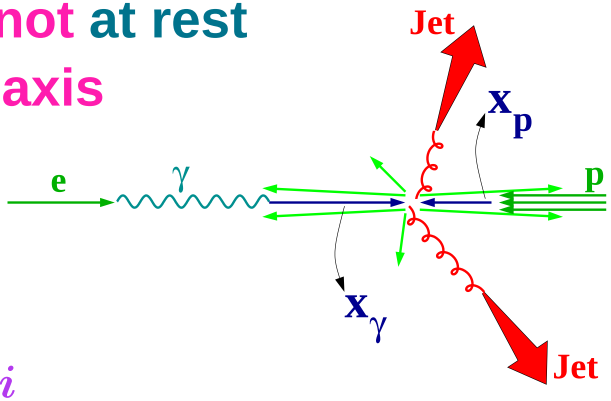


$\gamma p \rightarrow e + \text{jet} + \text{jet} + \text{X}$
(resolved photoproduction)

- Initial state: colored partons
- Initial partons carry only a fraction x_p, x_γ of the parent hadron
- Spectator partons:
  - $\longrightarrow$ remnant jets
  - $\longrightarrow$ "underlying event": soft interaction between the partons in the remnants

Jet search in ep collisions

- At HERA, ep collisions do not occur in the CM system:
  $\longrightarrow$ initial-state parton-parton or photon-parton system not at rest
  $\longrightarrow$ final-state partonic system boosted along the beam axis



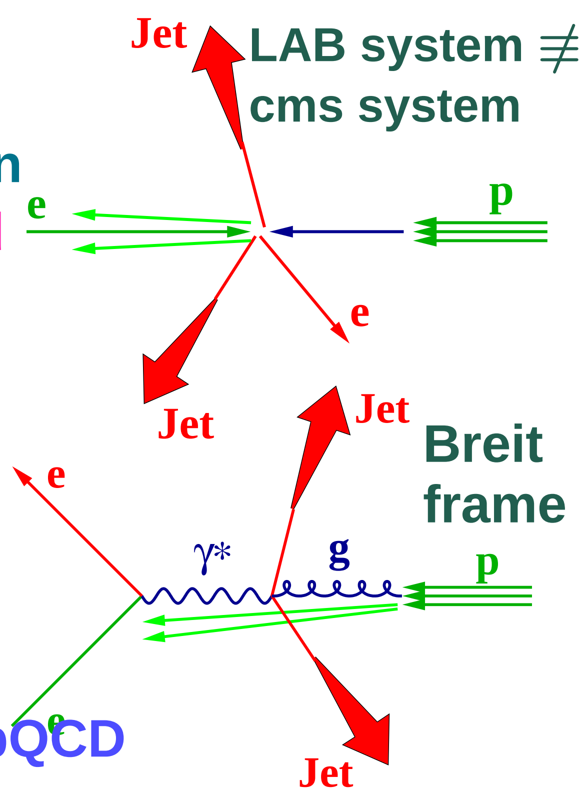
LAB system $\not\equiv$ centre-of-mass system

- To treat on equal footing all possible final-state hadronic systems $\longrightarrow$ use variables invariant under longitudinal boosts: $E_{T,i}$, $\Delta\eta_{ij}$ ($\eta = -\ln\tan\frac{\theta}{2}$) and φ_i

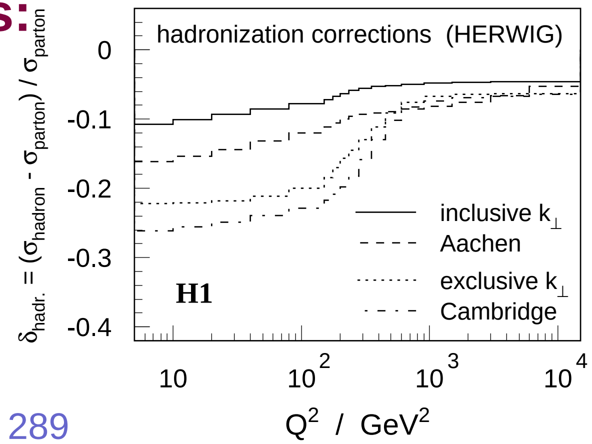
- Advantage of using transverse energies:
  $\longrightarrow$ large energy $\neq$ small distance (ie hard scattering)
  $\longrightarrow$ the beam remnant jets have huge energies, but they have not undergone a hard scattering
  $\longrightarrow$ large momentum transfer $\equiv$ small distance (hard scattering): large transverse energies $\equiv$ hard interaction
  $\longrightarrow$ the use of transverse energies helps to disentangle between the products of the hard interaction and the beam remnant jets (absent in e^+e^- annihilations)

Jet search in ep collisions

- **The kinematics of NC DIS poses several challenges:**
  - $\longrightarrow$ **presence of beam remnant jet**
  - $\longrightarrow$ **the initial-state γ^*-parton system is boosted (the parton carries a fraction of the proton momentum) and rotated (the γ^* carries p_T)**



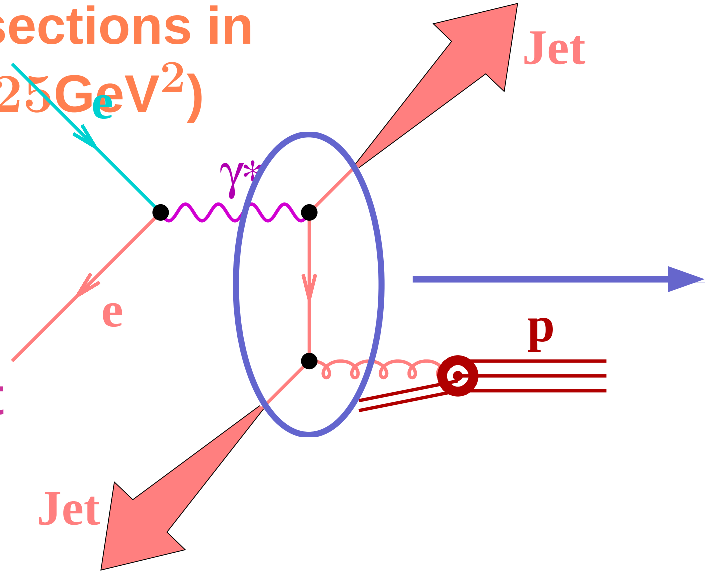
- **The effect of the p_T carried by the γ^* is removed by selecting a frame in which the γ^* collides head-on with the proton, eg the Breit frame**
- **k_T cluster algorithm in longitudinally inclusive mode:**
  - $\longrightarrow$ **best algorithm to reconstruct jets at HERA**
  - $\longrightarrow$ **in use since many years for making precision tests of pQCD**
- **Advantages of k_T algorithm in hadronic-type interactions:**
  - $\longrightarrow$ **allows transparent translation of experimental set-up to theoretical calculations (avoids ambiguities of overlapping and merging of jets (cone algorithm))**
  - $\longrightarrow$ **calculations using k_T are finite at all orders**
  - $\longrightarrow$ **smallest hadronisation corrections**



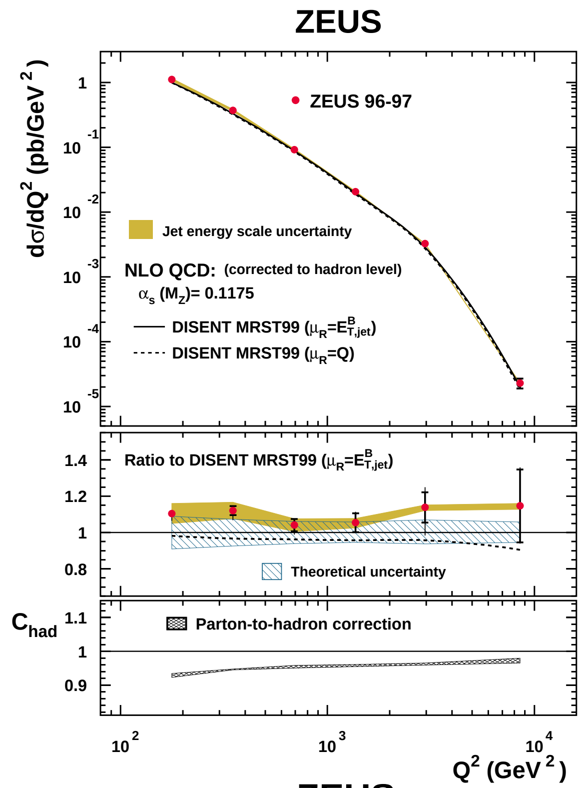
EPJ C19 (2001) 289

Inclusive-jet cross sections in NC DIS

- Measurements of jet cross sections in NC DIS at large Q^2 ($Q^2 > 125$ GeV^2) allow tests of pQCD and determination of α_s



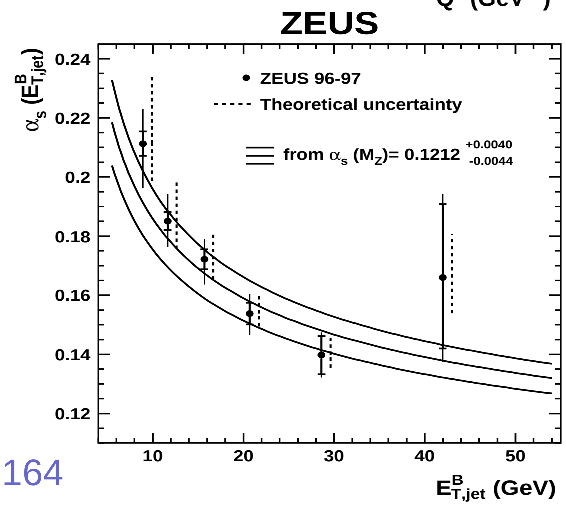
- The LO prediction for the jet cross section in the Breit frame is proportional to α_s
  $\rightarrow$ the measurements are directly sensitive to α_s:



$$\alpha_s(M_Z) = 0.1212 \pm 0.0017\ (\text{stat.})\ ^{+0.0023}_{-0.0031}\ (\text{exp.})\ ^{+0.0028}_{-0.0027}\ (\text{th.})$$

$\rightarrow$ experimental uncertainties: 2.9% ($Q^2 > 500$ GeV^2)
$\rightarrow$ theoretical uncertainties: 2.3%

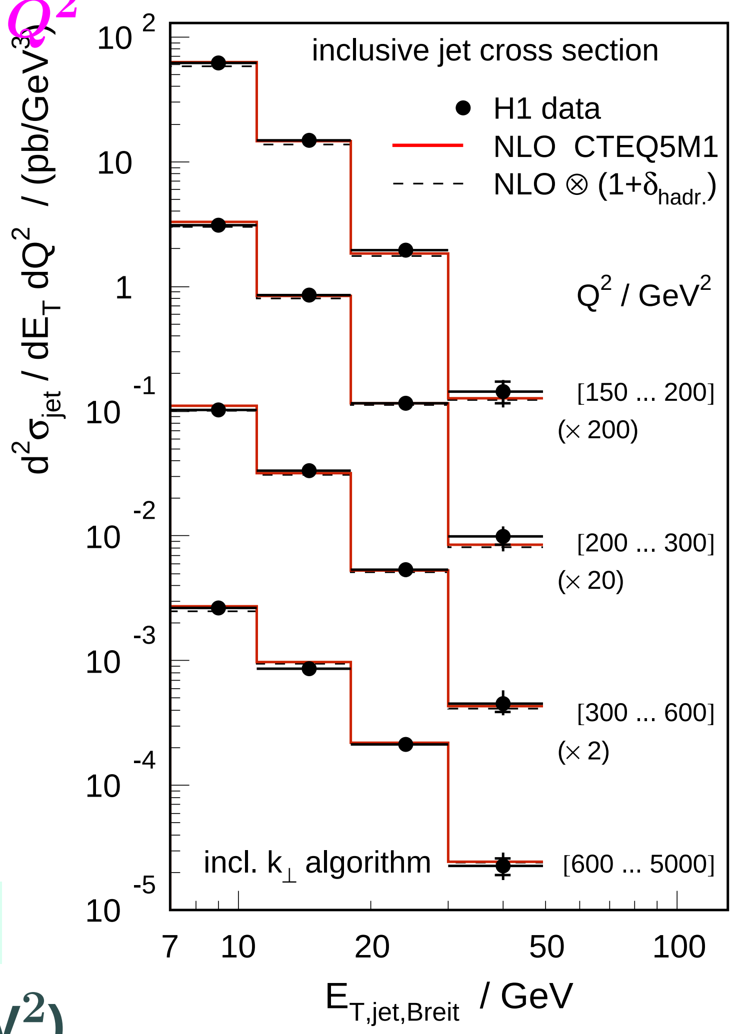
$\rightarrow$ Test of the energy-scale dependence of α_s from $d\sigma/dE^{\text{jet}}_{T,\text{B}}$ in NC DIS



PLB 547 (2002) 164

Inclusive-jet cross sections in NC DIS

- Measurements of the differential inclusive-jet cross section at large Q^2 ($Q^2 > 150$GeV^2) as a function of $E_{T,\mathrm{B}}^{\mathrm{jet}}$ in different Q^2 regions $\longrightarrow$ harder spectrum in $E_{T,\mathrm{B}}^{\mathrm{jet}}$ as Q^2 increases



- Comparison with NLO QCD calculations:
  - $\longrightarrow$ the predictions give a good description of the data over a wide range in Q^2 and $E_{T,\mathrm{B}}^{\mathrm{jet}}$
  - $\longrightarrow$ Validity of the description of the dynamics of inclusive jet production by pQCD at $\mathcal{O}(\alpha\alpha_s^2)$ over a wide range in Q^2 and $E_{T,\mathrm{B}}^{\mathrm{jet}}$

$\longrightarrow$ The measurements are directly sensitive to α_s:

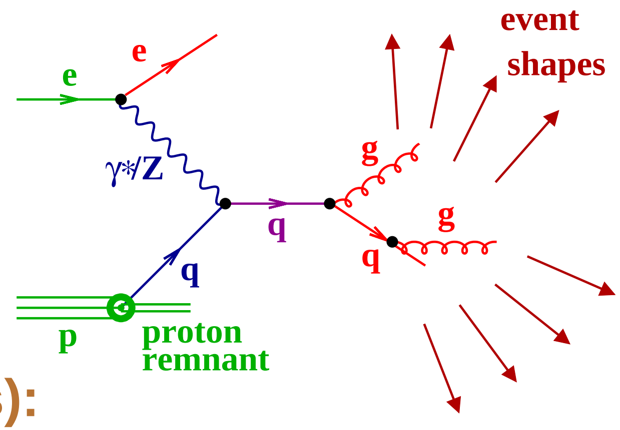
$$\alpha_s(M_Z) = 0.1186 \pm 0.0030\ (\mathrm{exp.}) \pm 0.0051\ (\mathrm{th.})$$

$\longrightarrow$ experimental uncertainties: 2.5% ($Q^2 > 150$ GeV^2)
$\longrightarrow$ theoretical uncertainties: 4.3%

EPJ C 19 (2001) 289

Hadronisation process

- The HFS in NC DIS has also been used to study the hadronisation process:
  $\longrightarrow$ non-perturbative effect
  $\longrightarrow$ event-shapes observables can be used to test the power-correction model: understanding of hadronisation process from first principles



- Event-shape variables (inspired by e^+e^- measurements):
  $\longrightarrow$ thrust, broadening, C parameter, jet mass

- In this type of analysis, the data are compared to a model prediction which consists of a combination of NLO QCD calculations and the expectations of the power corrections, characterised by an effective coupling $\bar{\alpha}_0$:

$$\boldsymbol{F} = \boldsymbol{F}_{\text{perturbative}} + \boldsymbol{F}_{\text{power correction}}$$

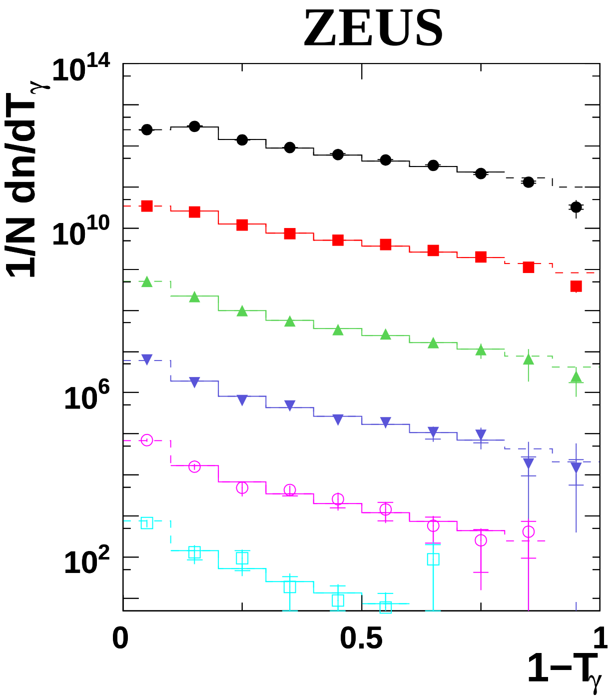
where $\boldsymbol{F}$ is an event-shape mean or distribution (NLO + matched NLL)

Test of the power-correction model

- **Extracted $\bar{\alpha}_0$ and $\alpha_s(M_Z)$ values for each event-shape observable:**

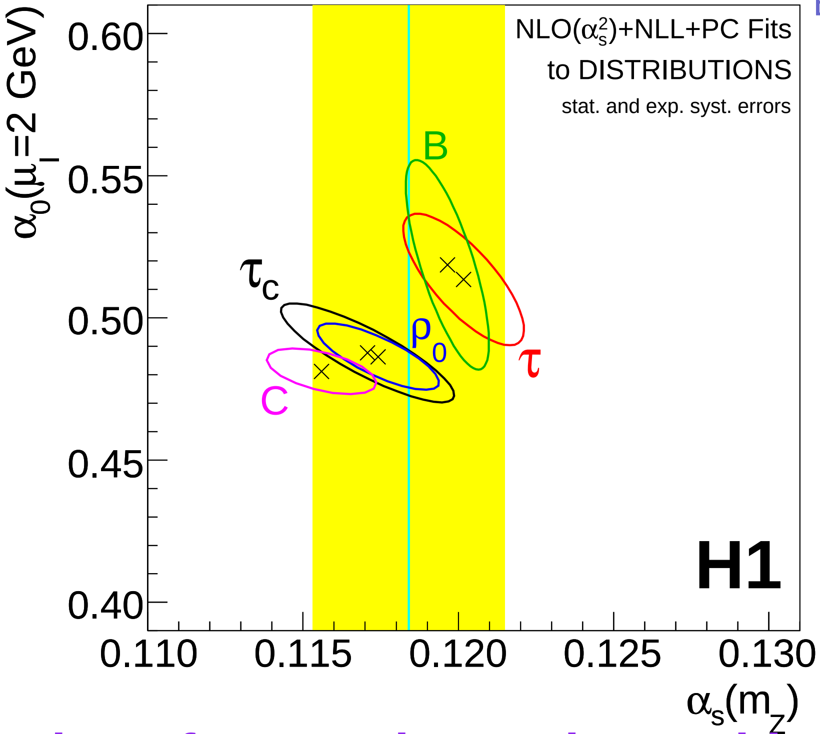
Thrust distribution in different Q^2 regions

DESY-06-042



EPJ C 46 (2006) 343

From distributions



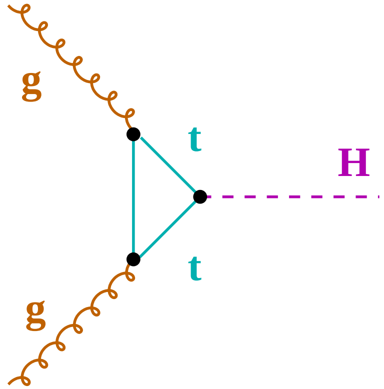
→ **It is possible to obtain good description of event-shape observables in hadronic-induced reactions using NLO+NLL+power corrections calculations**

→ **Extracted value of $\alpha_s(M_Z)$ consistent for all observables and with the world average**

→ **Universal non-perturbative parameter $\bar{\alpha}_0 = 0.5 \pm 10\%$**

→ **supports concept of power corrections as appropiate alternative approach for description of hadronisation effects**

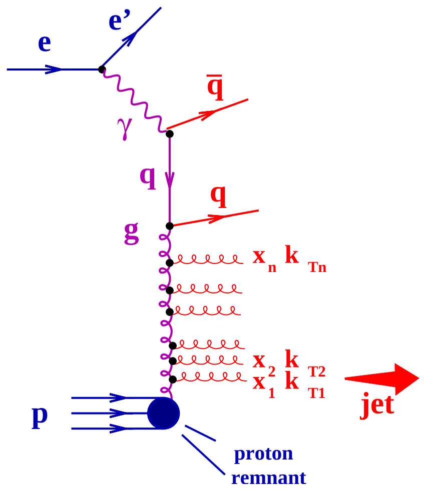
Parton evolution at low x and unintegrated PDFs



- One of the main channels of Higgs production at LHC is expected to be $gg \longrightarrow H$
- Predictions for these processes need information on
  - $\longrightarrow$ parton evolution at low x
  - $\longrightarrow$ unintegrated proton PDFs

  } $\longrightarrow$ forward-jet data at HERA
- At high scales (Q, $E_T^{\rm jet}$), calculations using the DGLAP evolution equations give a good description of the data at NLO

$\Rightarrow$ Measurements at HERA have provided
- $\longrightarrow$ accurate determination of the proton PDFs
- $\longrightarrow$ sensitive tests of pQCD and precise determinations of α_s



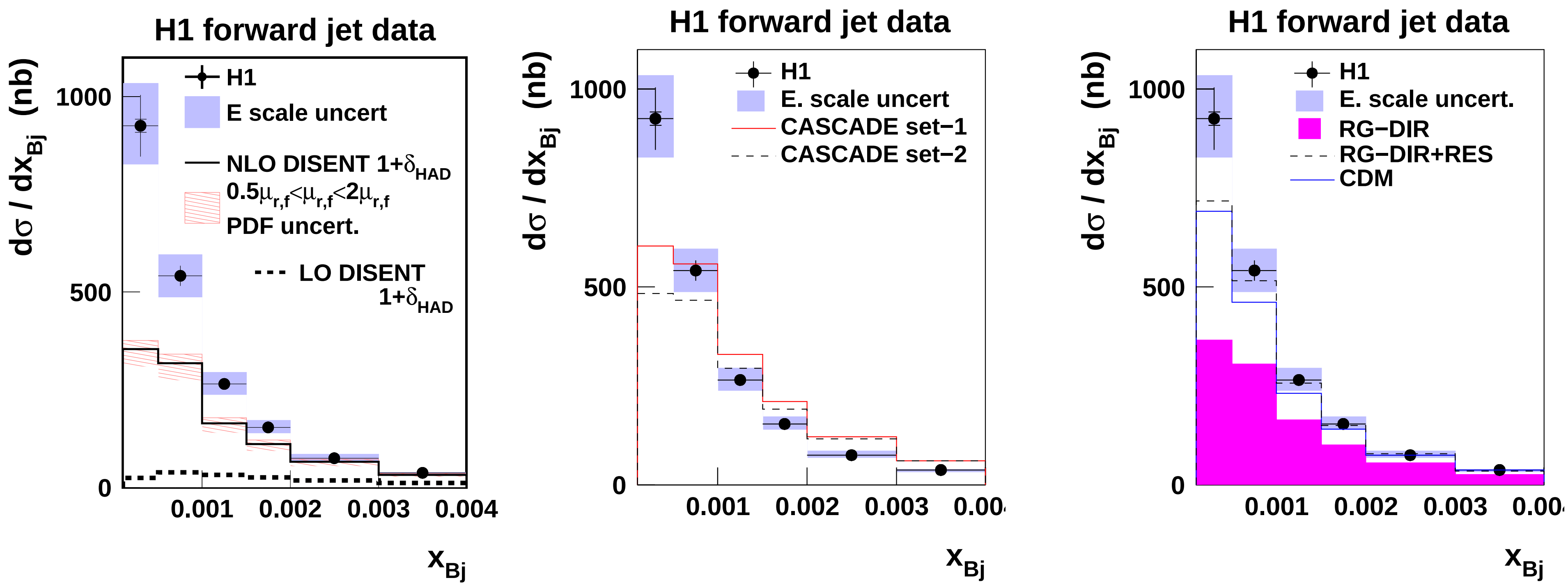
- DGLAP evolution equivalent to exchange of a parton cascade with exchanged partons strongly ordered in virtuality k_T
- But, DGLAP approximation expected to break down at low x:
  - — only leading logs in Q^2 are resummed
  - — contributions from $\log 1/x$ neglected (important for $\log Q^2 \ll \log 1/x$)

Parton evolution at low x and unintegrated PDFs

- HERA $\longrightarrow$ ideal testbed for theoretical approaches that account for low-x effects:
  - BFKL evolution: resummation of large $\log 1/x$ to all orders (very low x)
    $\longrightarrow$ no k_T ordering
    $\longrightarrow$ integration taken over full k_T phase space of gluons
    $\longrightarrow$ use of off-shell matrix elements together with unintegrated PDFs
  - CCFM evolution: angular-ordered parton emission (low and larger x)
    $\longrightarrow$ equivalent to BFKL for $x \rightarrow 0$ and to DGLAP at large x
    $\longrightarrow$ use of off-shell matrix elements together with unintegrated PDFs
  - virtual-photon structure: higher-order QCD effects mimicked at low x by introducing a second k_T-ordered parton cascade on the photon side
    $\longrightarrow$ resolved is expected to contribute for $(E_T^{\rm jet})^2 > Q^2$ and suppressed with increasing Q^2
- By restricting jet data to
  $\longrightarrow$ large $\eta^{\rm jet}$ (forward direction, proton side)
  $\longrightarrow$ $x_{\rm jet} = E^{\rm jet}/E_p \gg x_{\rm Bj}$ (to suppress QPM)
  $\longrightarrow$ $Q \approx E_T^{\rm jet}$ (to restrict evolution in Q^2)
  these different approaches has been investigated

Forward jet production at low x in NC DIS

EPJ C46 (2006) 27



⟶ Forward-jet cross section rises with decreasing x

- Comparison to predictions:
  - ⟶ NLO pQCD (DGLAP): fails to describe the data at low x
  - ⟶ CASCADE (CCFM): improved description of data at low x (sensitivity to unintegrated PDFs)
  - ⟶ CDM (BFKL-like) and resolved-photon: better description of data

Azimuthal jet separation

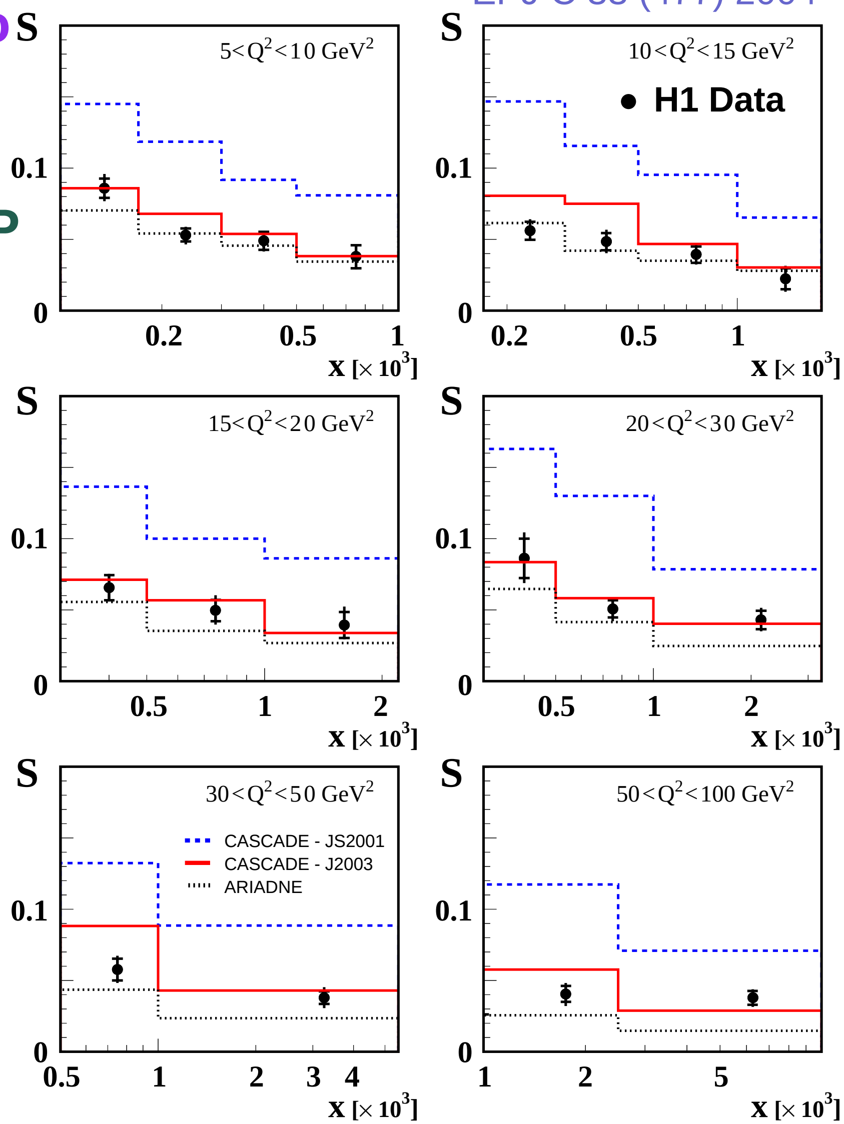
EPJ C 33 (477) 2004

- Insight into low-x dynamics can be gained also by studying the azimuthal separation between the two hardest jets: an excess of events at small $\Delta\phi$ would signal a deviation from DGLAP evolution

- The ratio

$$\longrightarrow \quad S = \frac{\int_0^{\alpha} N_{2\text{jet}}(\Delta\phi^*, x, Q^2) d\Delta\phi^*}{\int_0^{\pi} N_{2\text{jet}}(\Delta\phi^*, x, Q^2) d\Delta\phi^*}, \quad \alpha = \frac{2}{3}\pi$$

is well suited to test small-x effects



- Comparison to CCFM predictions (CASCADE):
  $\longrightarrow$ calculations of S show sensitivity to the unintegrated gluon distributions

$\longrightarrow$ These measurements can be used to constrain the unintegrated PDFs

Dijet cross sections in photoproduction

- Measurements of jet cross sections in photoproduction allow tests of color dynamics
- At HERA, quark and gluon exchange can be studied in the same hadronic-induced reaction by separating resolved and direct processes in PHP using

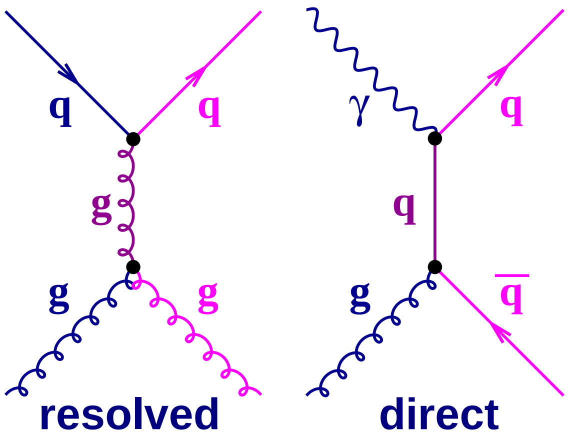
$$x_\gamma^{\rm obs} = \frac{1}{2yE_e}(E_T^{\rm jet1} e^{-\eta^{\rm jet1}} + E_T^{\rm jet2} e^{-\eta^{\rm jet2}}) \quad \rightarrow \begin{cases} x_\gamma^{\rm obs} < 1 \text{ for resolved} \\ x_\gamma^{\rm obs} \sim 1 \text{ for direct} \end{cases}$$

- Resolved processes dominated by gluon exchange (like dijets in pp):

$$\theta^* \rightarrow 0, \pi: \ \frac{d\sigma}{d\cos\theta^*} \sim \frac{1}{(1 - |\cos\theta^*|)^2}$$

- Direct processes proceed via quark exchange (like prompt photon in pp):

$$\theta^* \rightarrow 0: \ \frac{d\sigma}{d\cos\theta^*} \sim \frac{1}{(1 - |\cos\theta^*|)^1}$$



→ The $\cos\theta^*$ distribution reflects the underlying parton dynamics since it has sensitivity to the spin of the exchanged particle in two-body processes

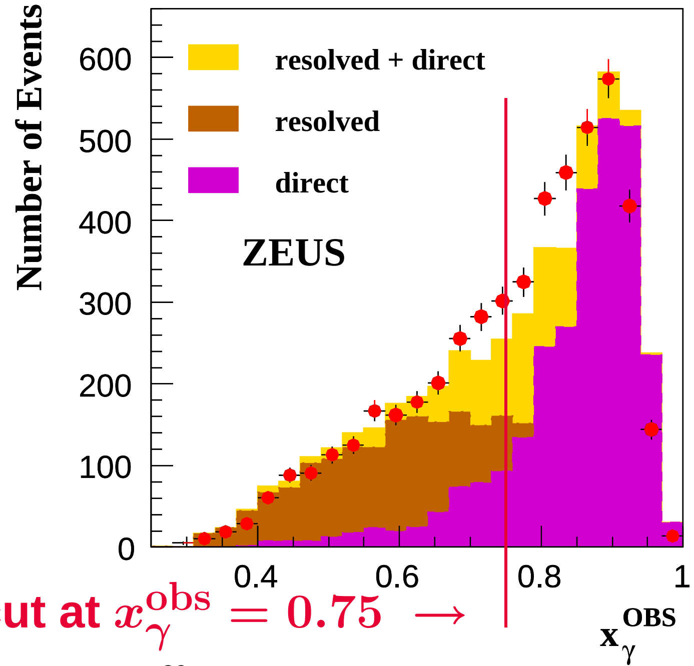
Dijet cross sections in photoproduction

- Measurements of jet cross sections in photoproduction allow tests of color dynamics

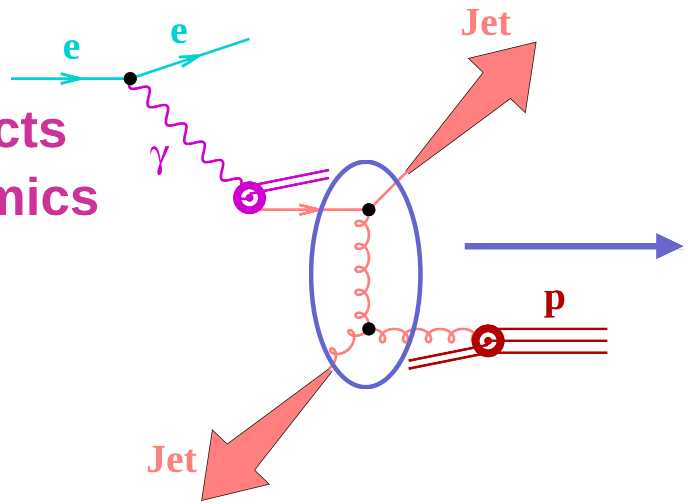
sample with $x_\gamma^{\rm obs} < 0.75 \leftarrow$ dominated by resolved processes

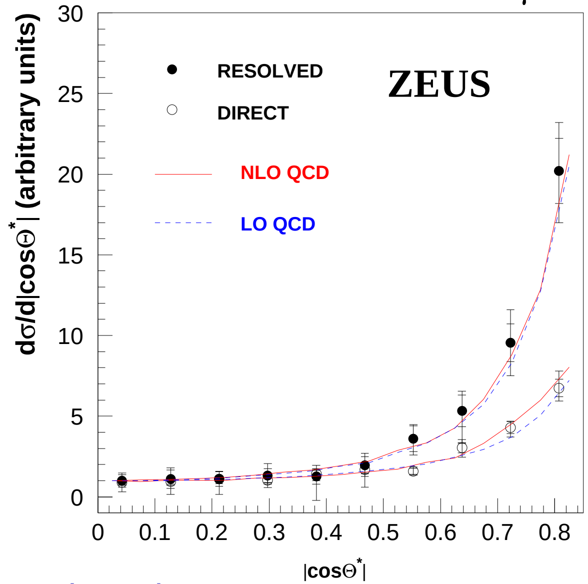
sample with $x_\gamma^{\rm obs} > 0.75 \leftarrow$ dominated by direct processes

cut at $x_\gamma^{\rm obs} = 0.75 \rightarrow$



- The $\cos\theta^*$ distribution reflects the underlying parton dynamics

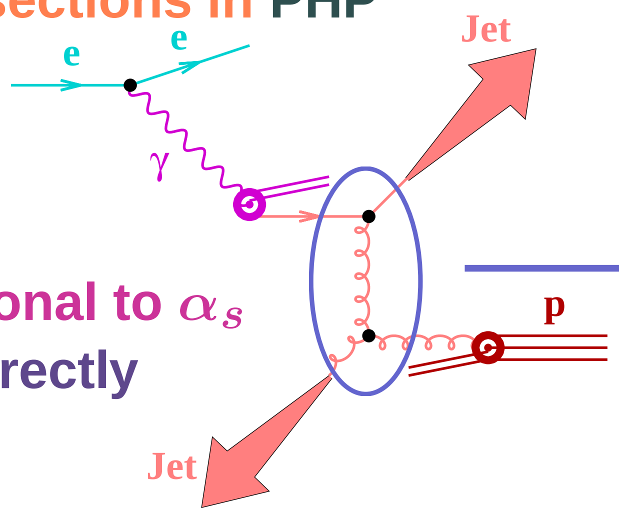


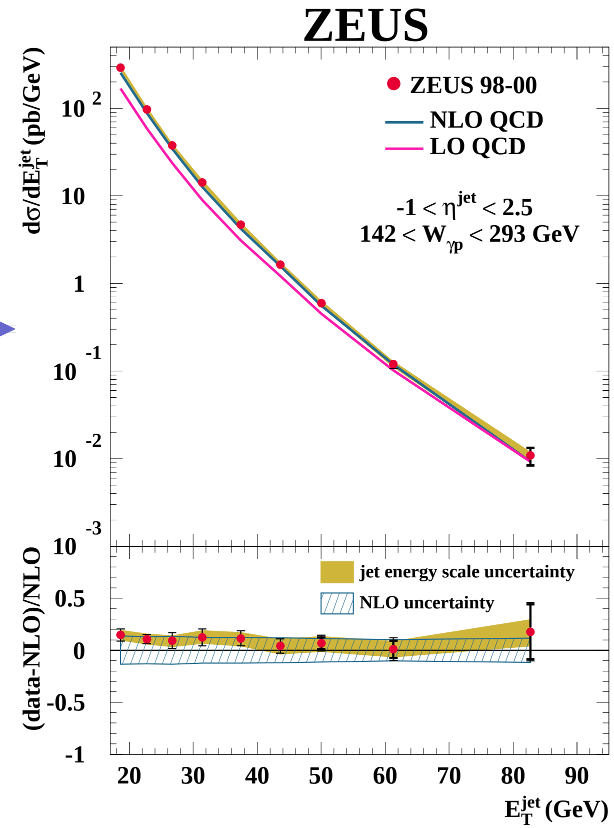


PLB 384 (1996) 401

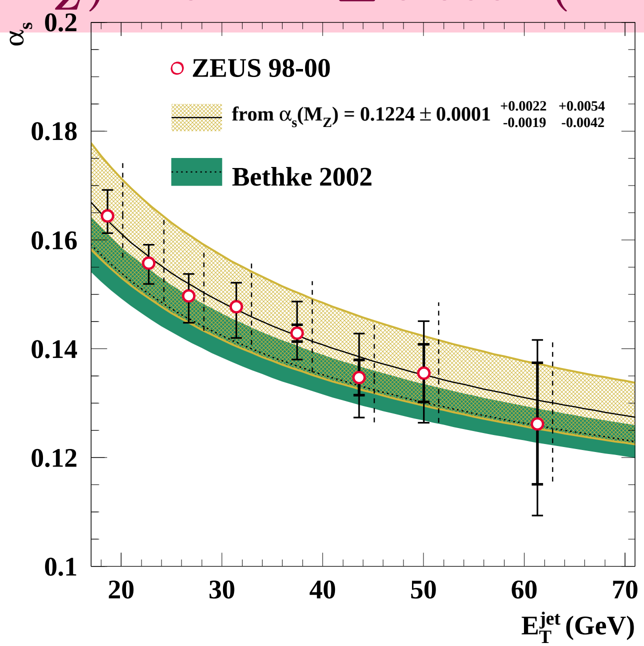
Inclusive-jet cross sections in photoproduction

- Measurements of jet cross sections in PHP at large E_T^{jet} allow tests of pQCD
- The LO prediction for the jet cross section is proportional to α_s
  $\longrightarrow$ the measurements are directly sensitive to α_s:





$$\alpha_s(M_Z) = 0.1224 \pm 0.0001 \text{ (stat.)} {}^{+0.0022}_{-0.0019} \text{ (exp.)} {}^{+0.0054}_{-0.0042} \text{ (th.)}$$



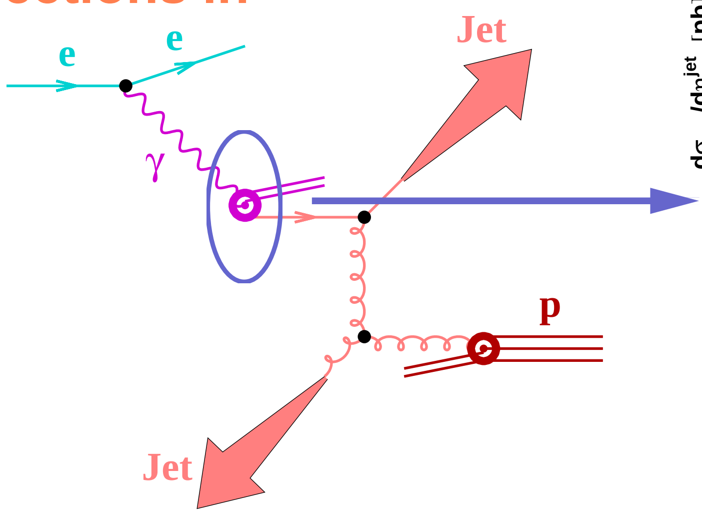
experimental uncertainties: 1.8%
theoretical uncertainties: 4.4%

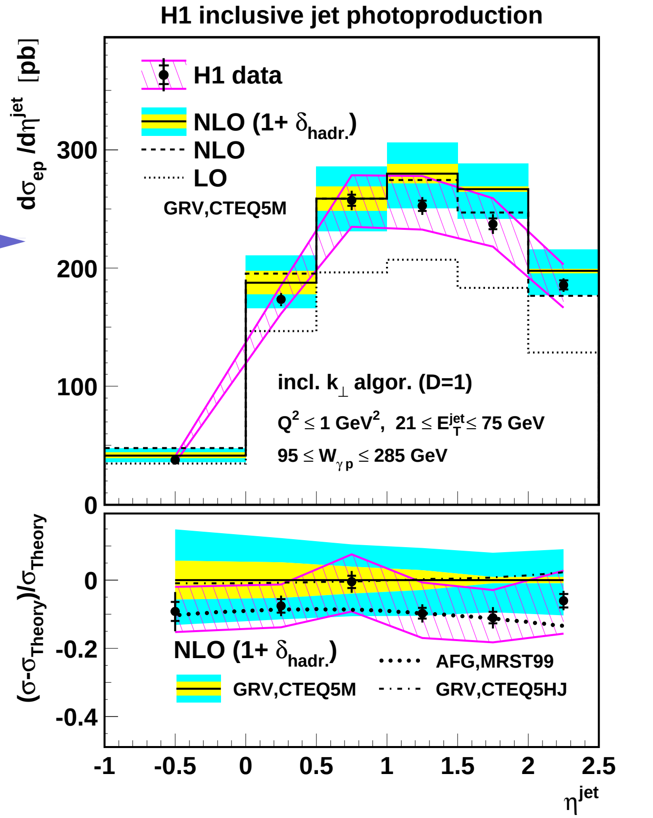
$\longrightarrow$ Test of the energy-scale dependence of α_s from $d\sigma/dE_T^{jet}$ in PHP

PLB 560 (2003) 7

Inclusive-jet cross sections in photoproduction

- Measurements of jet cross sections in photoproduction allow tests of γPDFs





EPJ C 29 (2003) 497

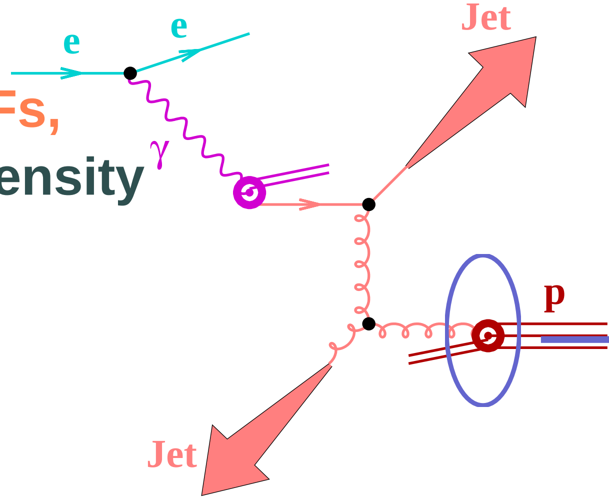
$$x_\gamma^{\rm obs} = \frac{1}{2yE_e}(E_T^{\rm jet1}e^{-\eta^{\rm jet1}} + E_T^{\rm jet2}e^{-\eta^{\rm jet2}})$$

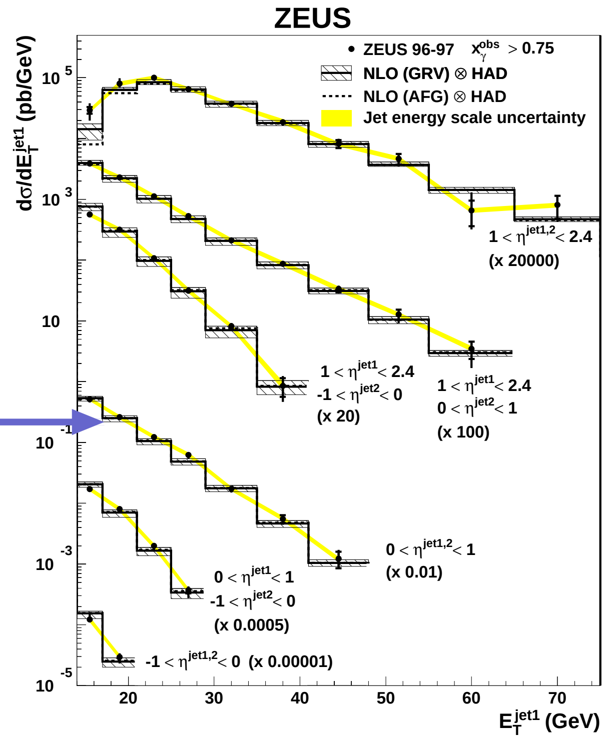
($x_\gamma^{\rm obs} < 1$ for resolved and $x_\gamma^{\rm obs} \sim 1$ for direct)

- The structure of the photon is investigated by measuring jet cross sections most sensitive to the γPDF's, eg $d\sigma/d\eta^{\rm jet}$ or $d\sigma/dx_\gamma^{\rm obs}$, and comparing the measurements to predictions based on different parametrisations of the γPDF's
  $\longrightarrow$ the measurements can be used to discriminate among different parametrisations or used in a global fit to constrain them

Dijet cross sections in photoproduction

- Measurements of jet cross sections in **photoproduction** provide a useful constrain of the pPDFs, in particular, on the **gluon density**





EPJ C 23 (2002) 4

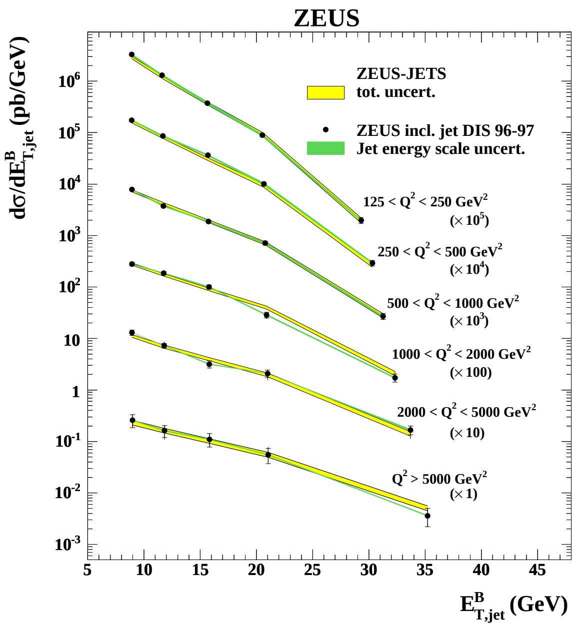
- This is achieved by measuring jet cross sections less sensitive to the γPDF's, eg $d\sigma/dE_T^{\rm jet}$ for $x_\gamma^{\rm obs} > 0.75$

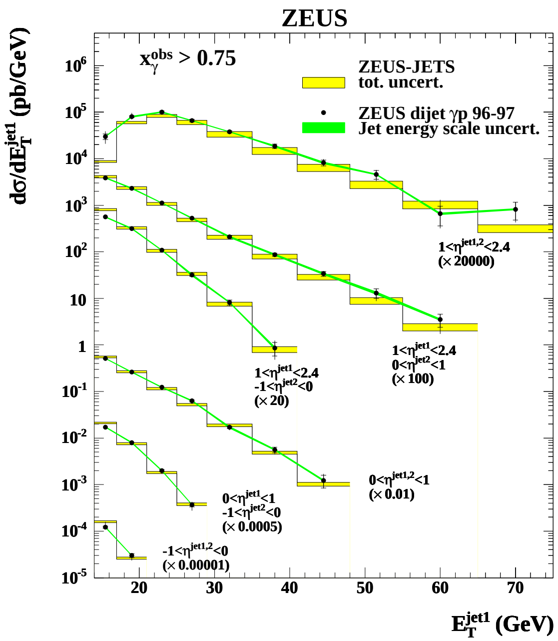
- These measurements have been incorporated in a global fit of the pPDFs to improve the determination of the gluon density in the proton

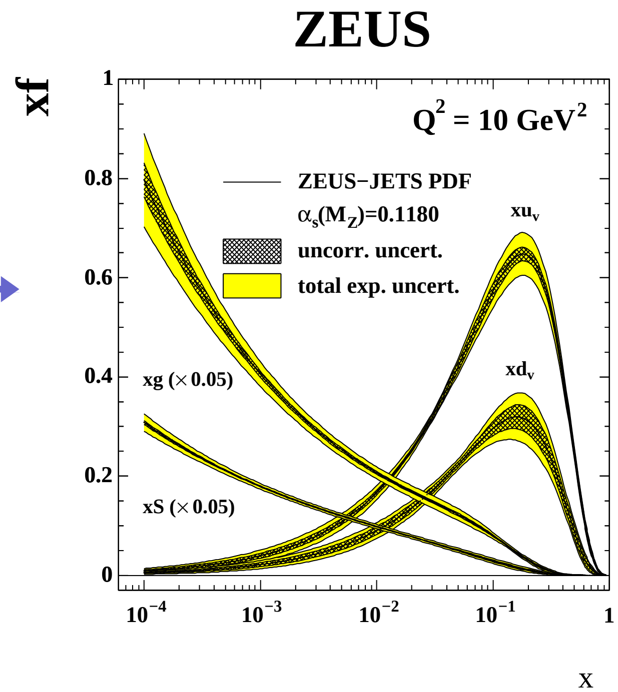
EPJ C 42 (2005) 1

Conclusions: jets and PDFs

- Very precise jet cross sections in NC DIS and PHP that are directly sensitive to the gluon content of the proton: useful to constrain gluon density, especially at mid- to high-x (most relevant at LHC energies)
- Measurements incorporated in a QCD fit to determine PDFs parametrisations:



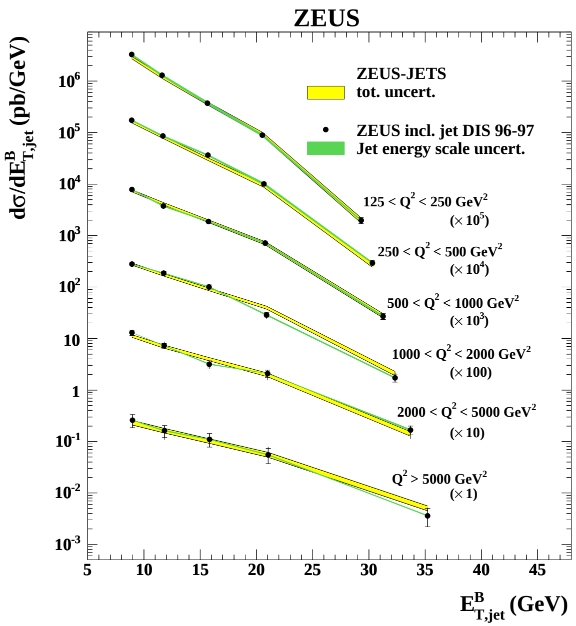


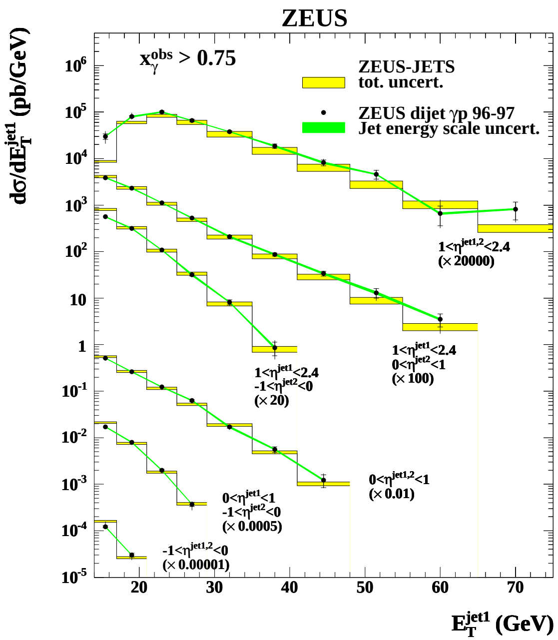


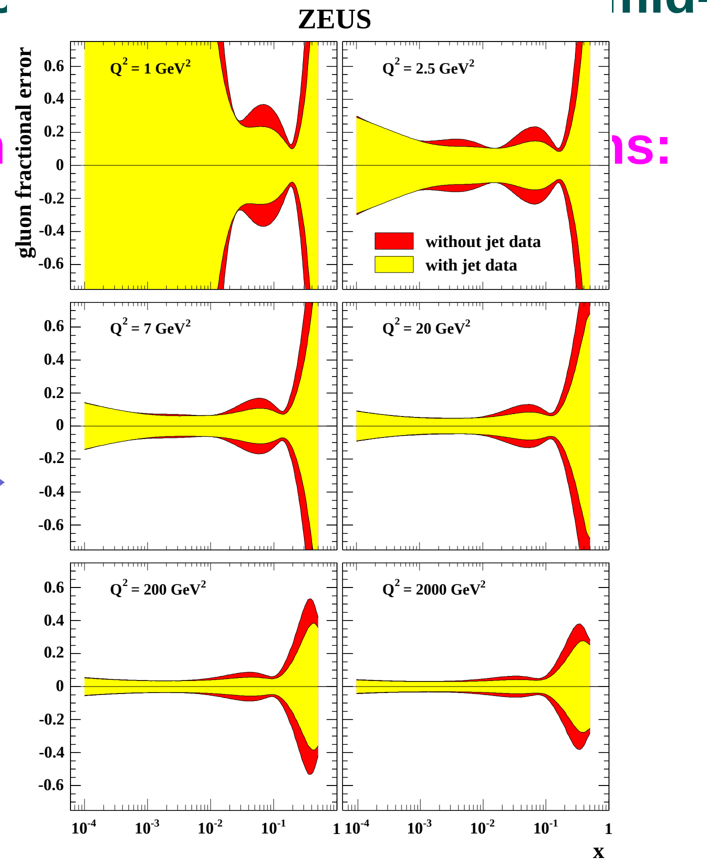
EPJ C 42 (2005) 1

Conclusions: jets and PDFs

- Very precise jet cross sections in NC DIS and PHP are directly sensitive to the gluon content of the proton: useful to constrain gluon density, especially at mid- to high-x (most relevant at LHC energies)
- Measurements incorporated in a QCD fit to determin ... ns:







- The result is an improvement on the determination of the gluon density in the proton → the uncertainty in the gluon density decreases for mid- to high-x by up to a factor of 2

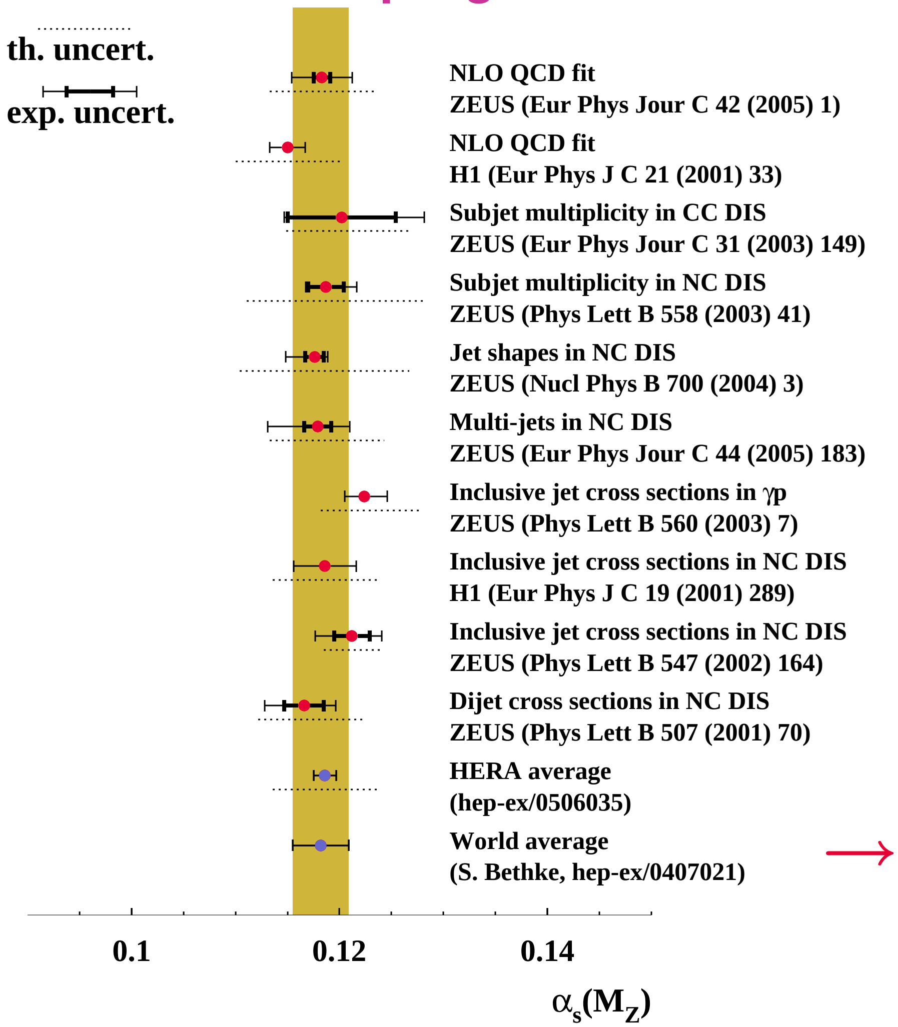
EPJ C 42 (2005) 1

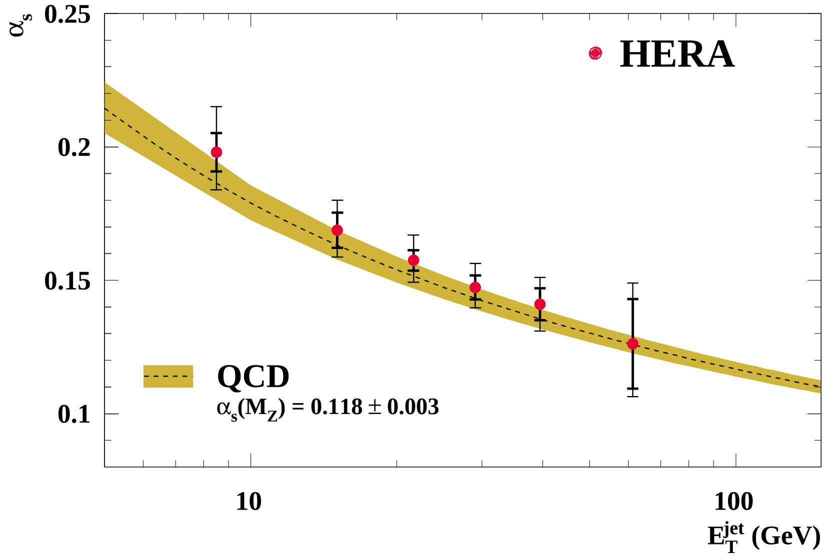
Conclusions: α_s

- **HERA has become a unique QCD-testing machine: very useful for understanding multijet production in "clean" hadronic-induced reactions**

  → considerable progress in understanding and reducing uncertainties led to very precise determinations of α_s

  ⇒ improved calculations needed for better accuracy





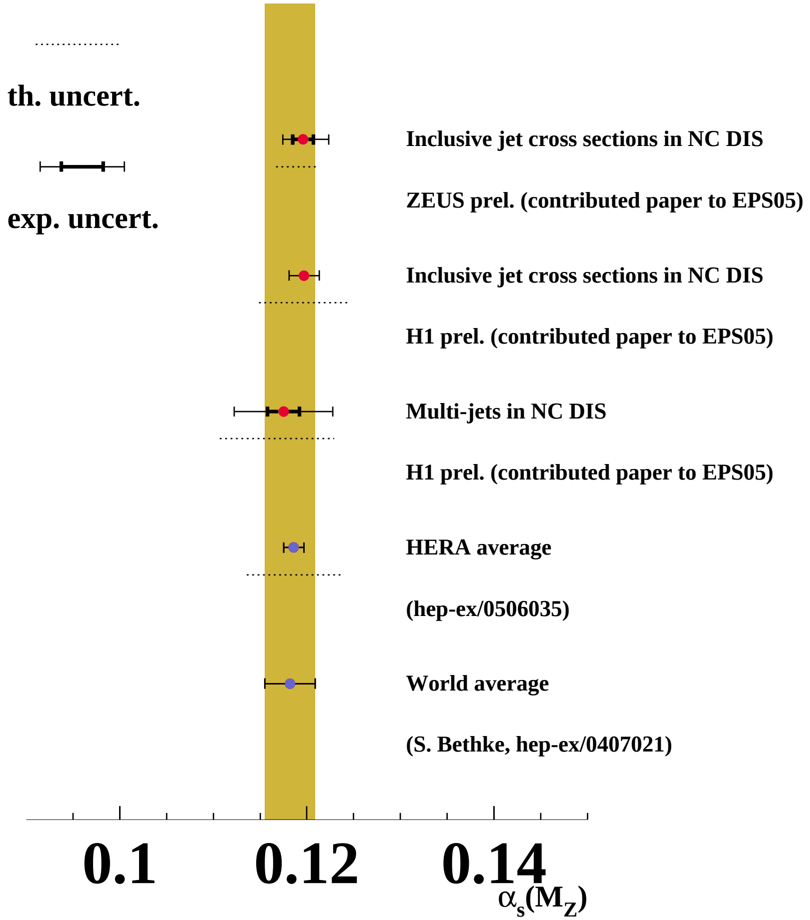
→ **Observation of the running of α_s from HERA jet data alone**

hep-ex/0506035

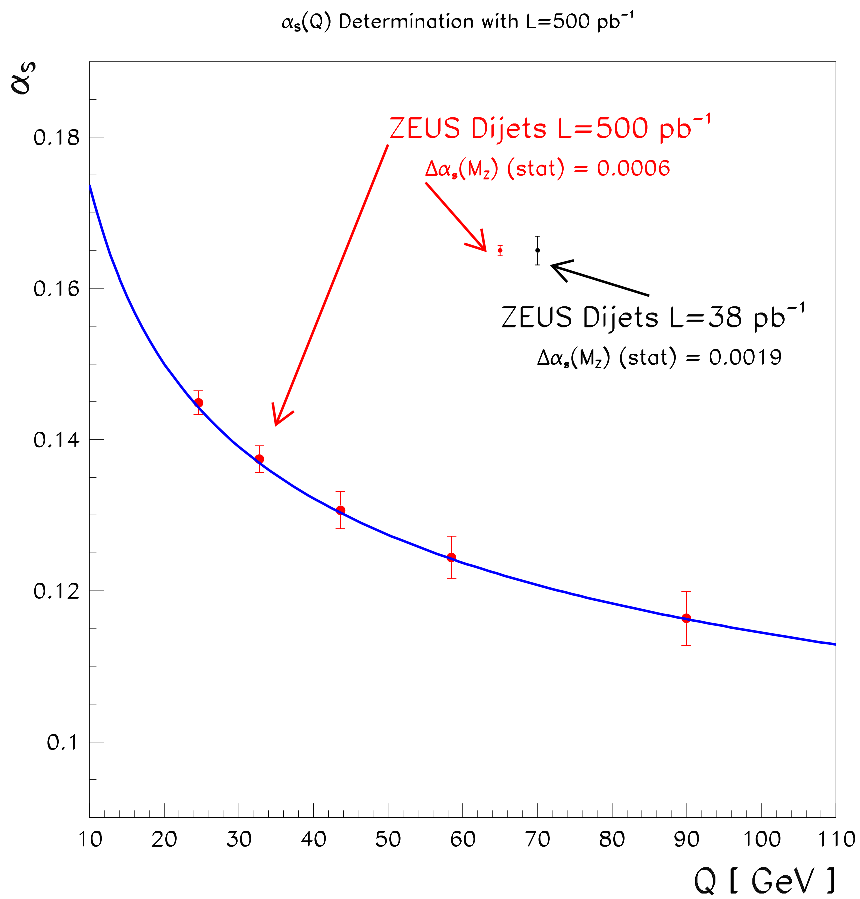
The future of HERA

- **HERA is still running: wealth of data already available to test new MC models, MC@NLO and NNLO pQCD calculations if they are made available for ep collisions**

$\longrightarrow$ **new HERA I results still coming:**



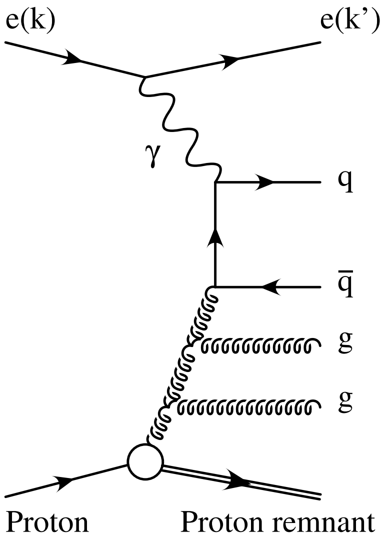
$\longrightarrow$ **85 pb^{-1} e^+p and 300 pb^{-1} e^-p HERA II data being analysed:**

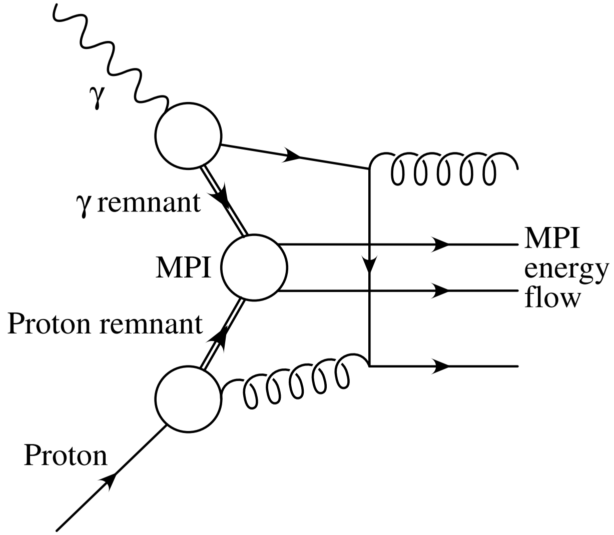


Back-up slides

Multijet cross sections in photoproduction

- Measurements of jet cross sections in photoproduction allow tests of hard multijet production, multiparticle interactions and "underlying event"

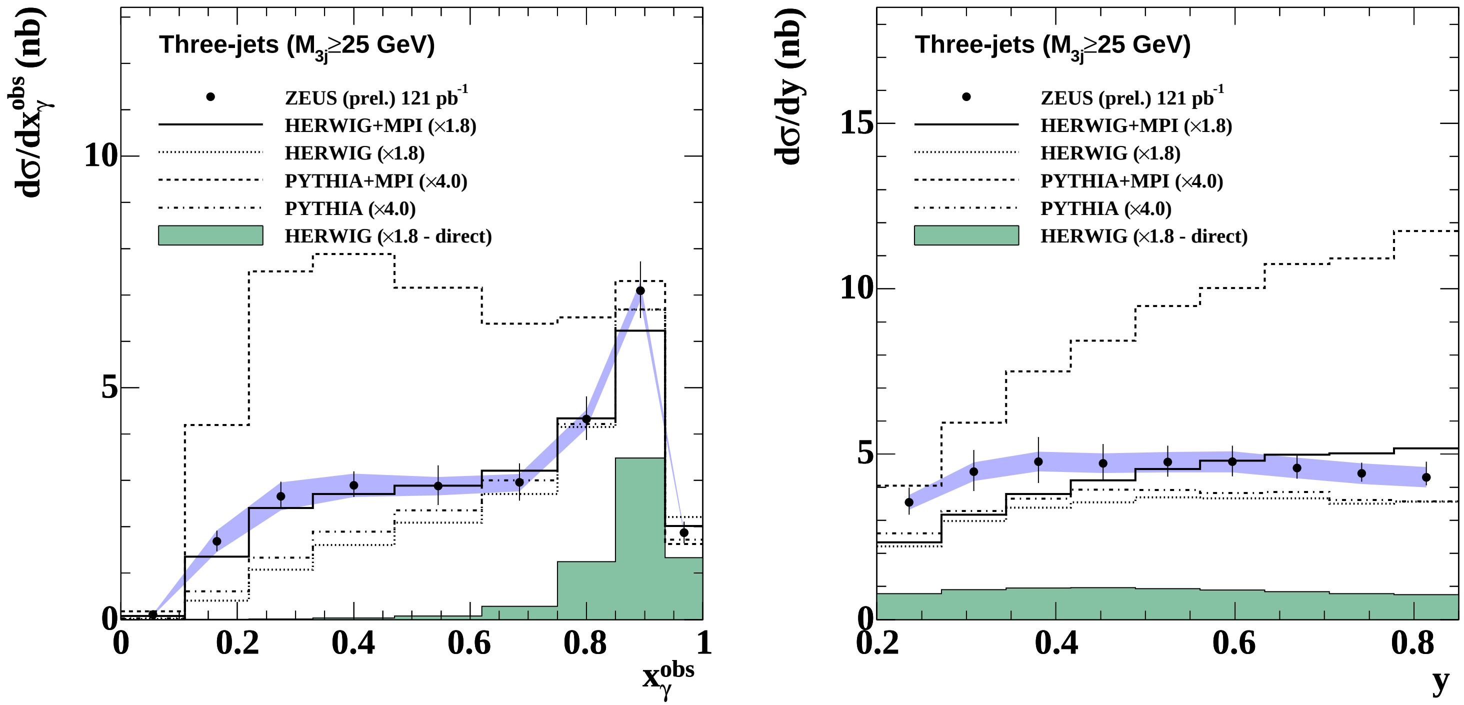




- Multijet production is directly sensitive to high orders: $\sigma_{3\text{jet}} \propto \alpha_s^2$
- Test of parton showers in Monte Carlo models
- Sensitivity to multiparton interactions/underlying event $\longrightarrow$ test/tune models

Multijet cross sections in photoproduction

DIS06 preliminary



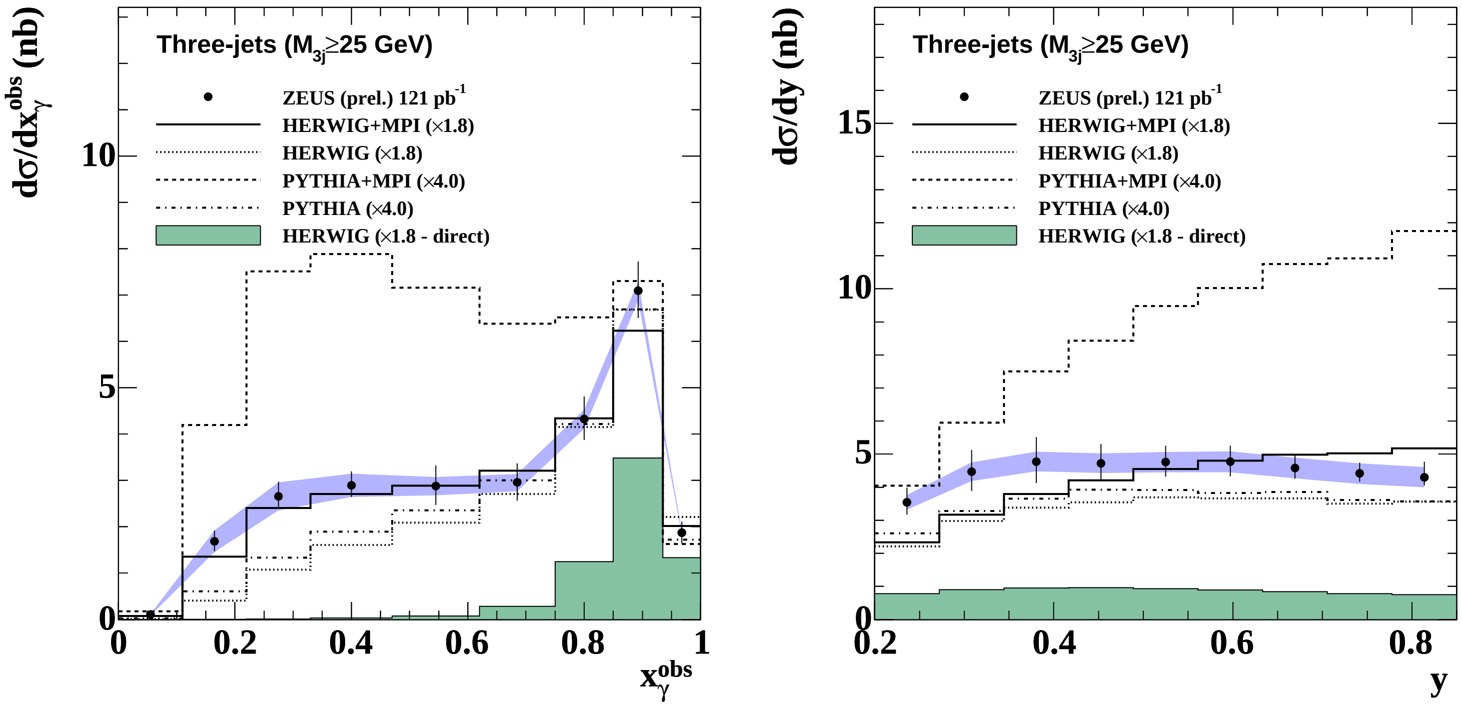
PYTHIA MPI tuned to generic collider data **HERWIG MPI tuned to $x_\gamma^{\rm obs}$ data**

- **Monte Carlo models without multiparton interactions describe the y ($=\sqrt{Ws}$) but fail describe the shape of the measured $x_\gamma^{\rm obs}$ distribution**
- **PYTHIA MPI fails to describe the data**
- **HERWIG MPI describes the $x_\gamma^{\rm obs}$ but the description of y gets spoiled**

$\rightarrow$ **$x_\gamma^{\rm obs}$ and y distributions: ideal ground for tuning and testing models**

Multijet cross sections in photoproduction

DIS06 preliminary



→ **very precise hadronic data can be used as testing ground for hadronisation and multiparton interaction models**